

Influence of spin-orbit coupling on the transport properties of spintronics materials

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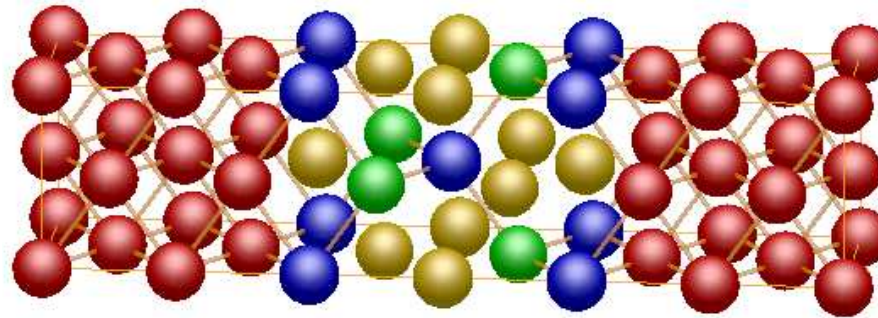
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Univ. Regensburg

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- Introduction
 - electronic structure calculations
- Transport in trilayer systems
 - TMR
 - in/out of plane anisotropy TAMR
 - in plane TAMR
- Magnetotransport in bulk
 - formalism
 - Residual resistivity tensor
 - spin transport
- Summary

FM/SC/FM-trilayers with perfect matching



Green's function (GF)

$$G^{\pm}(\vec{r}, \vec{r}'; E) = \lim_{\epsilon \rightarrow 0} \sum_{\lambda} \frac{\phi_{\lambda}(\vec{r}) \phi_{\lambda}^{\times}(\vec{r}')}{E - E_{\lambda} \pm i\epsilon} = G^{\pm}(\vec{r}_i + \vec{R}_i, \vec{r}'_j + \vec{R}_j, E)$$

tight-binding version of KKR-method

$$G^{+}(\vec{r}, \vec{r}'; E) = \frac{1}{A_{\text{SBZ}}} \int_{\text{SBZ}} d^2 k_{\parallel} e^{i\vec{k}_{\parallel}(\vec{\rho}_{\nu} - \vec{\rho}_{\nu'})}$$

$$\times \left[\sum_{\Lambda \Lambda'} R_{\Lambda}^{\nu}(\vec{r}, E) G_{\Lambda \Lambda'}^{\nu \nu'}(\vec{k}_{\parallel}, E) R_{\Lambda'}^{\nu'}(\vec{r}', E) - i p \delta_{\nu \nu'} \sum_{\Lambda} R_{\Lambda}^{\nu}(\vec{r}_{<}, E) H_{\Lambda}^{\nu'}(\vec{r}_{>}, E) \right]$$

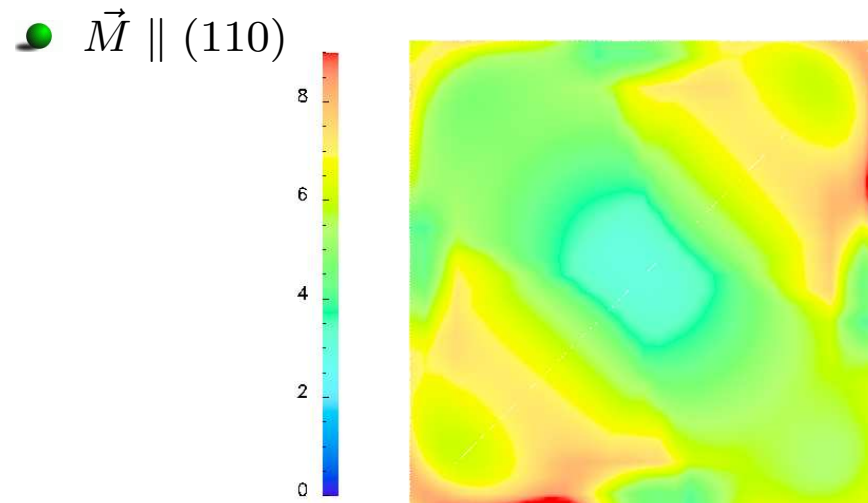
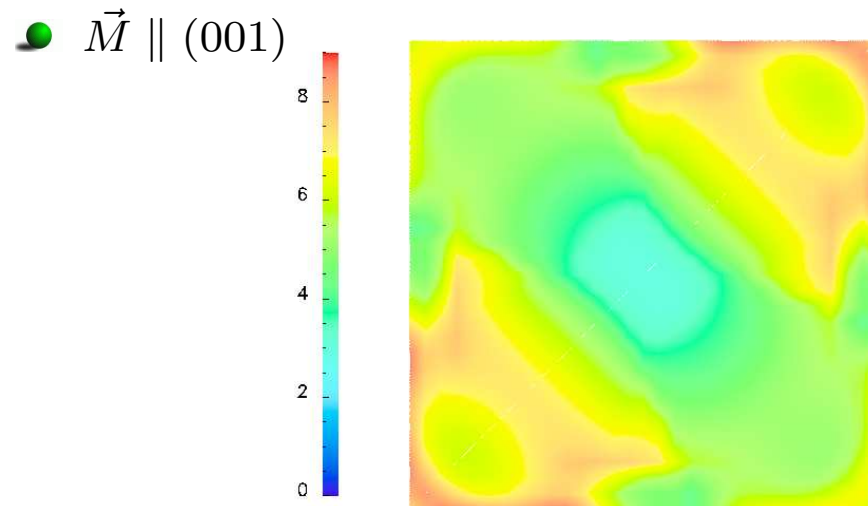
Dirac Hamiltonian within LSDA (Local Spin Density Approximation)

$$\begin{aligned}
 \hat{H}_D &= c\vec{\alpha}\vec{p} + \beta mc^2 + V(\vec{r}) + \beta B(\vec{r})\sigma_z, \quad \vec{\alpha} = \begin{pmatrix} 0 & \vec{\sigma} \\ \vec{\sigma} & 0 \end{pmatrix}, \quad \beta = \begin{pmatrix} I_2 & 0 \\ 0 & -I_2 \end{pmatrix} \\
 &= \left[i\gamma_5 \sigma_r c \left(\frac{\partial}{\partial r} + \frac{1}{r} (1 - \beta \hat{K}) \right) + V(r) + \beta \sigma_z B(r) + (\beta - 1) \frac{c^2}{2} \right] \\
 \hat{K} &= \hat{\sigma} \hat{L} + 1, \quad \gamma_5 = \begin{pmatrix} 0 & -I_2 \\ -I_2 & 0 \end{pmatrix}, \quad \sigma_r = \frac{1}{r} \vec{r} \cdot \vec{\sigma}
 \end{aligned}$$

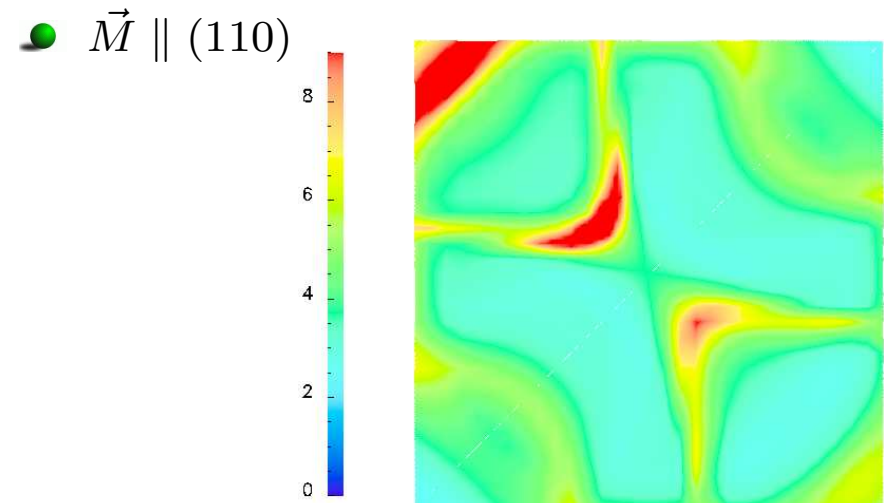
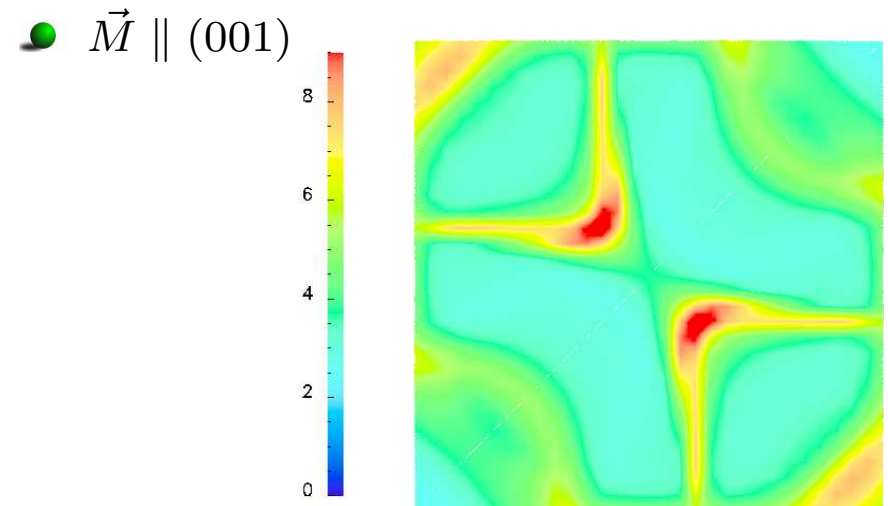
- four component Dirac formalism
- accounts for SOC and spin-polarisation on same level

Bloch spectral functions $A(\vec{k}_{\parallel}, E_F)$ – SC interface layer

As-termination

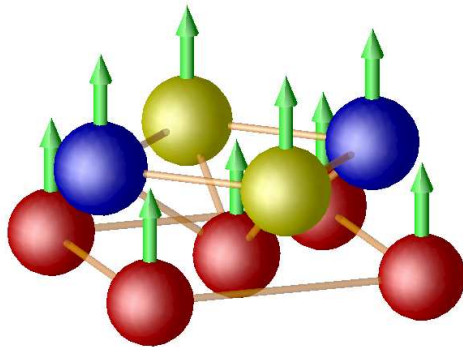


Ga-termination

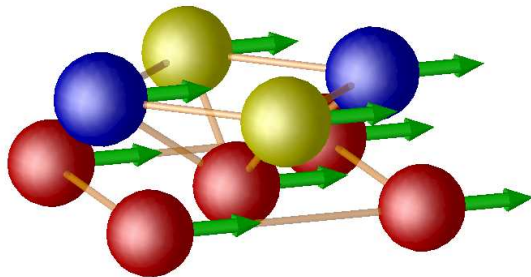


Symmetry breaking at the FM/SC-interface

● $\vec{M} \parallel (001) - C_{2v}$



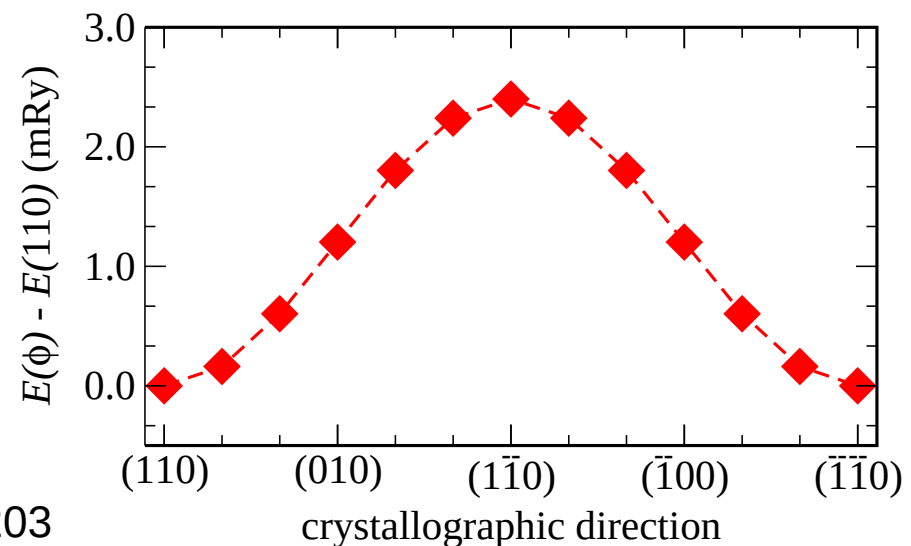
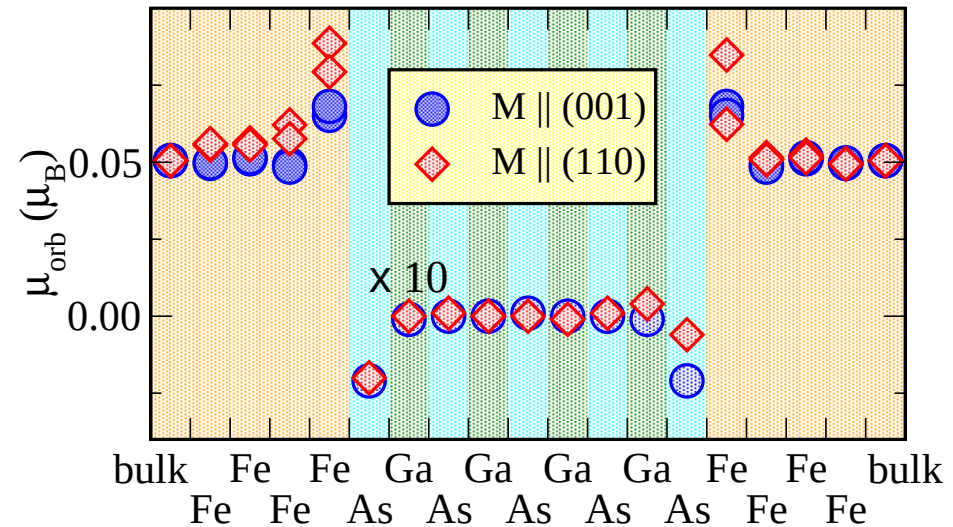
● $\vec{M} \parallel (110) - C_s$



see also: Sjöstedt *et al.* 2002, PRL **89**, 267203

Košuth *et al.* 2005, EPL **72**, 816

● Fe/9(GaAs)/Fe - μ_{orb} and MAE



Conductance - Landauer-Büttiker formalism

$$G = \frac{e^2}{h} \sum_{\vec{k}_{\parallel}} g(\vec{k}_{\parallel})$$

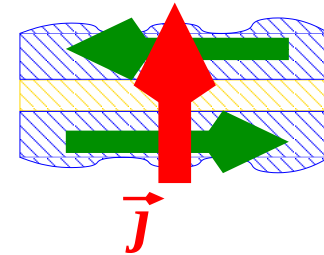
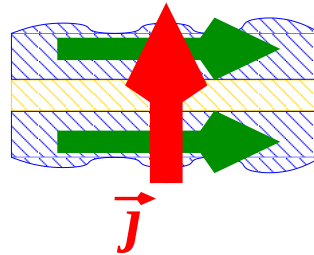
with $\vec{k}_{\parallel}$ -resolved conductance:

$$g(\vec{k}_{\parallel}) = \int_{A_{\text{WS}}^{\text{L}}} d^2r \int_{A_{\text{WS}}^{\text{R}}} d^2r' j_z(\vec{r}) G(\vec{r}, \vec{r}'; \vec{k}_{\parallel}; E_{\text{F}}) j_z(\vec{r}') G^*(\vec{r}, \vec{r}'; \vec{k}_{\parallel}; E_{\text{F}})$$

$$j_z(\vec{r}) = \frac{c}{E_{\text{F}} + mc^2} \left(-i\hbar \nabla_z + \frac{V}{c} \vec{\alpha}_z - i\frac{B}{c} \beta (\vec{\alpha} \times \vec{a}_z)_z \right)$$

”pessimistic” MR ratio

$$T = \frac{G^{\text{P}} - G^{\text{AP}}}{G^{\text{P}}}$$



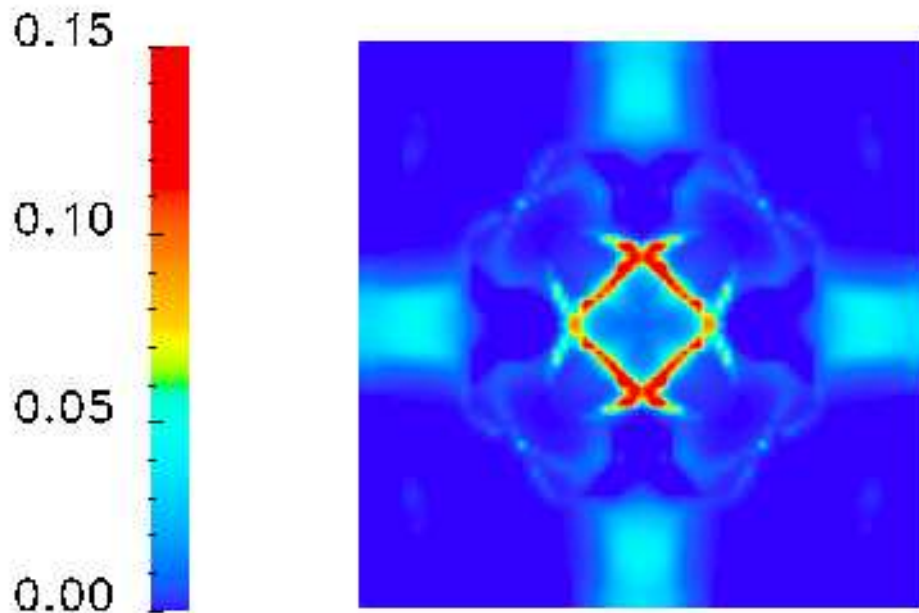
P / AP : parallel / antiparallel orientation of magnetisation KKR:

Mavropoulos *et al.* (2004)

$\vec{k}_{||}$ -resolved conductance – AP orientation

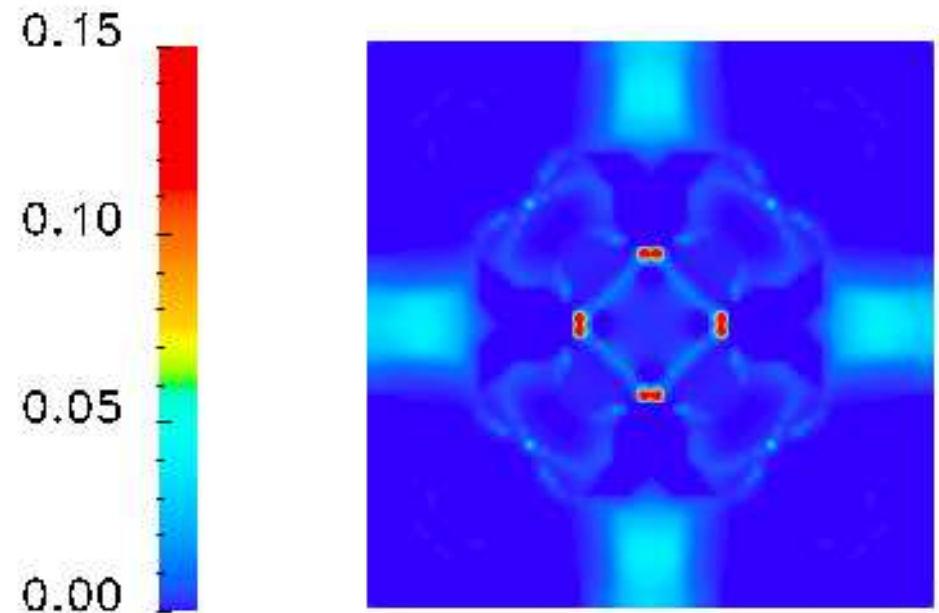
- effect of the spin-orbit coupling (SOC)

correct SOC



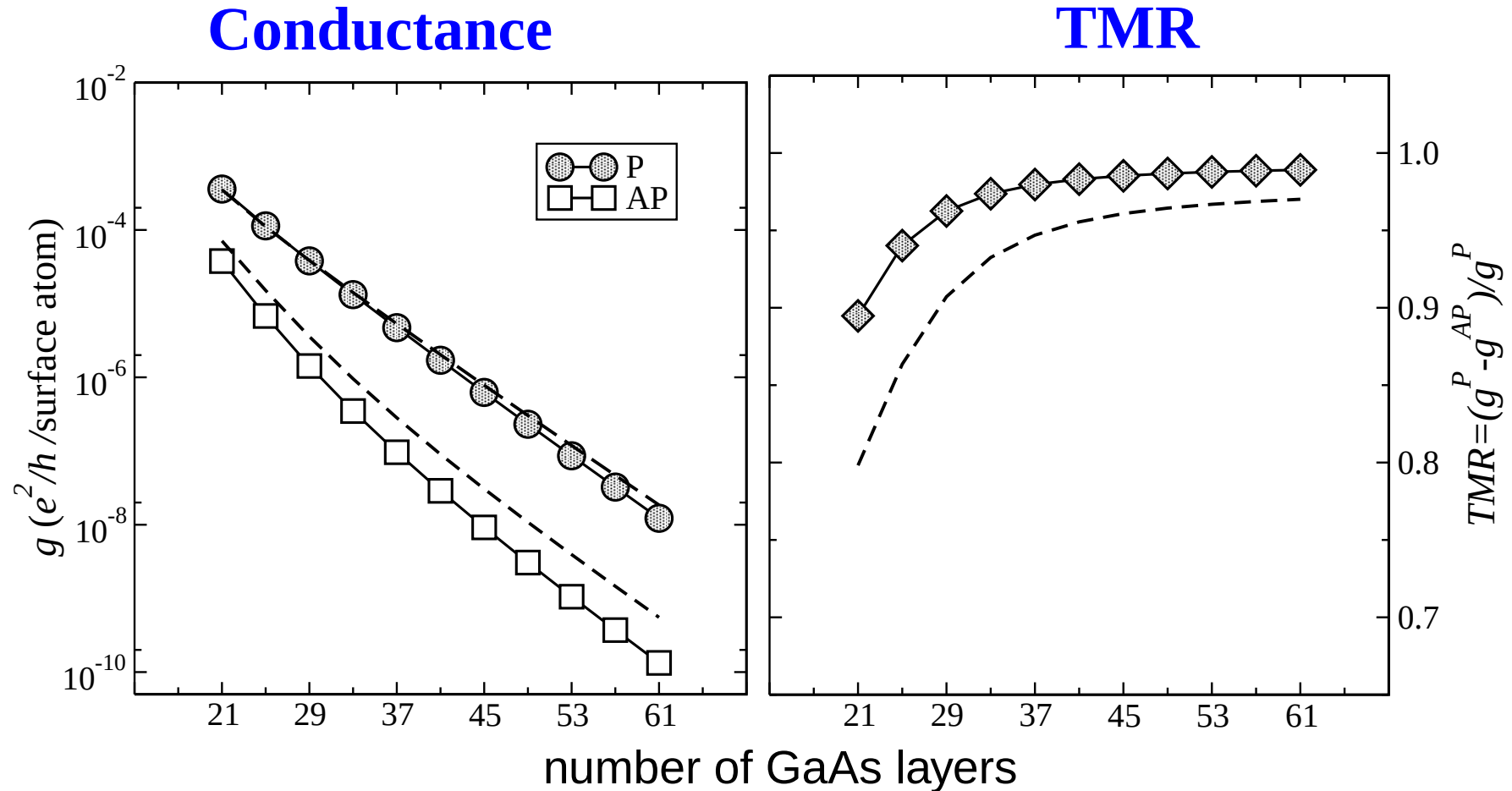
$$G^{\text{AP}} = 0.0076 \, e^2/h$$

SOC switched off



$$G^{\text{AP}} = 0.0044 \, e^2/h$$

Tunneling Conductance for $\text{Fe}/n(\text{GaAs})/\text{Fe}$

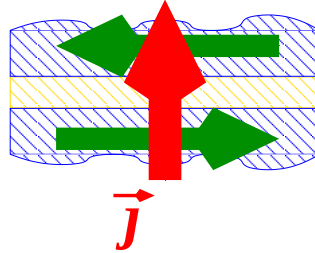


- Symbols: SOC suppressed for inner GaAs layers
- dashed lines: full SOC $\Rightarrow$ SOC increases g^{AP}

Dependence on orientation of $\vec{M}$

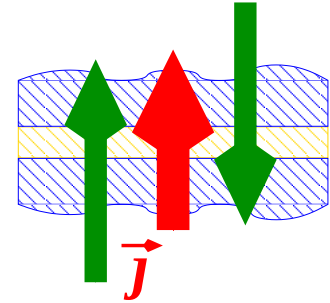
$\vec{M}$ along (110)

● AP orientation

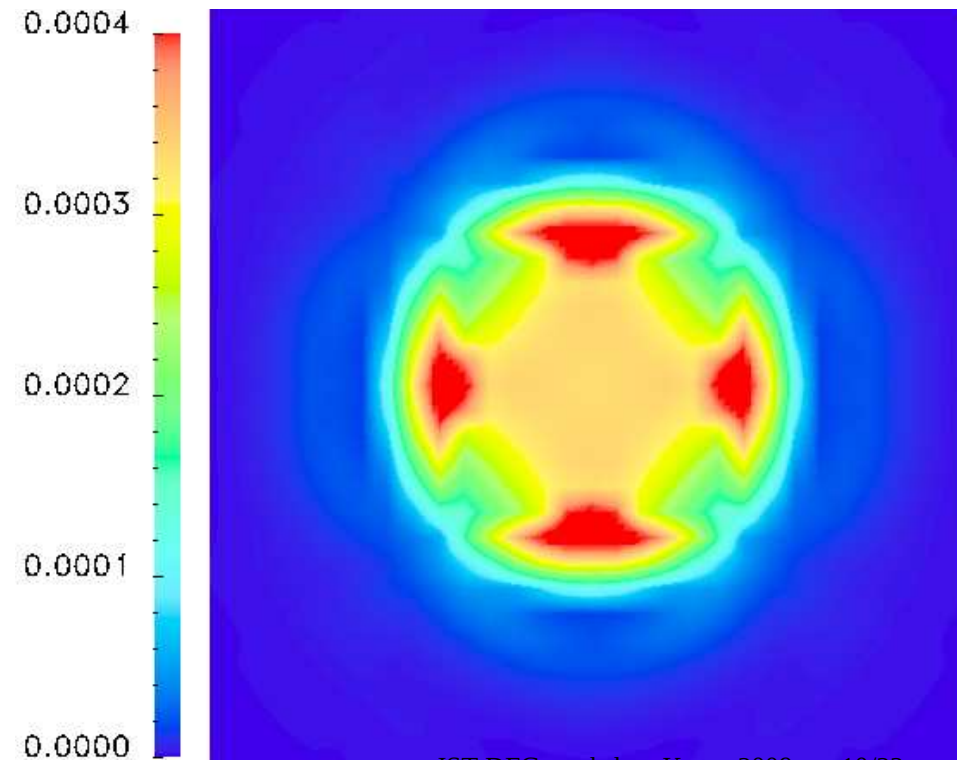
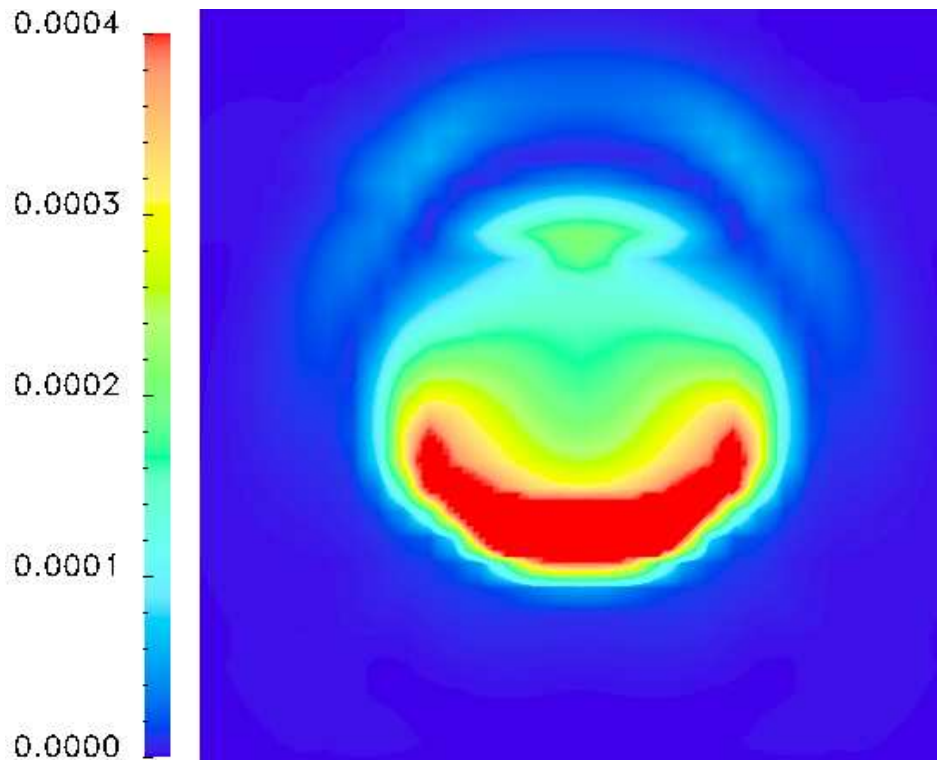


$\vec{M}$ along (001)

● AP orientation



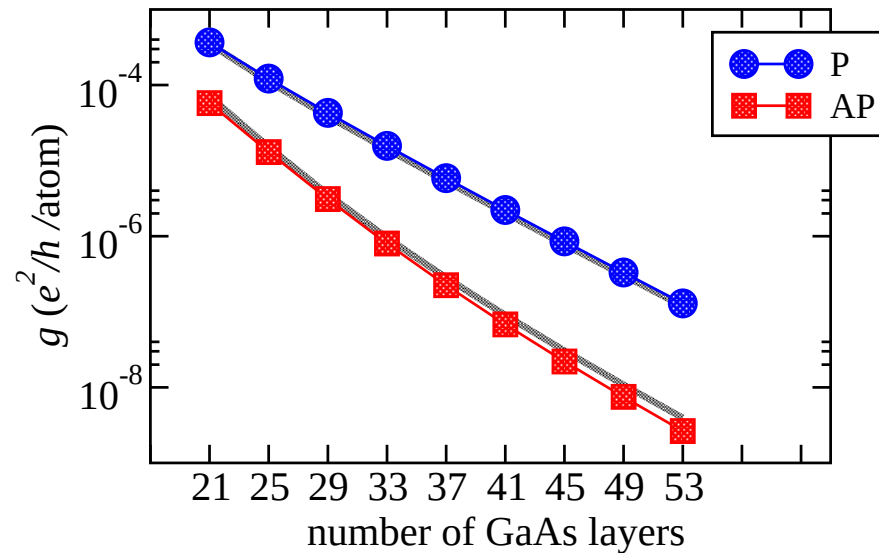
$k_{||}$ -resolved conductance



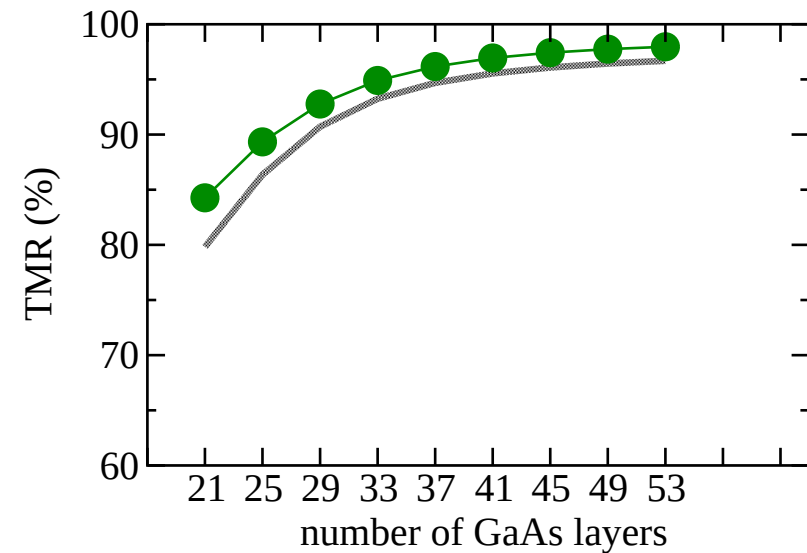
Dependence on orientation of $\vec{M}$

Conductance for $\vec{M} \parallel (110)$ and (001)

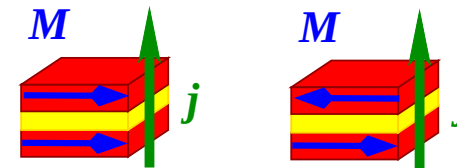
● conductance for P/AP alignment



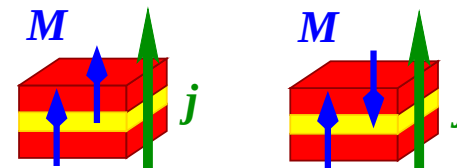
● $\text{TMR} = (G^P - G^{AP})/G^P$



● symbols for $\vec{M} \parallel (110)$



lines for $\vec{M} \parallel (001)$

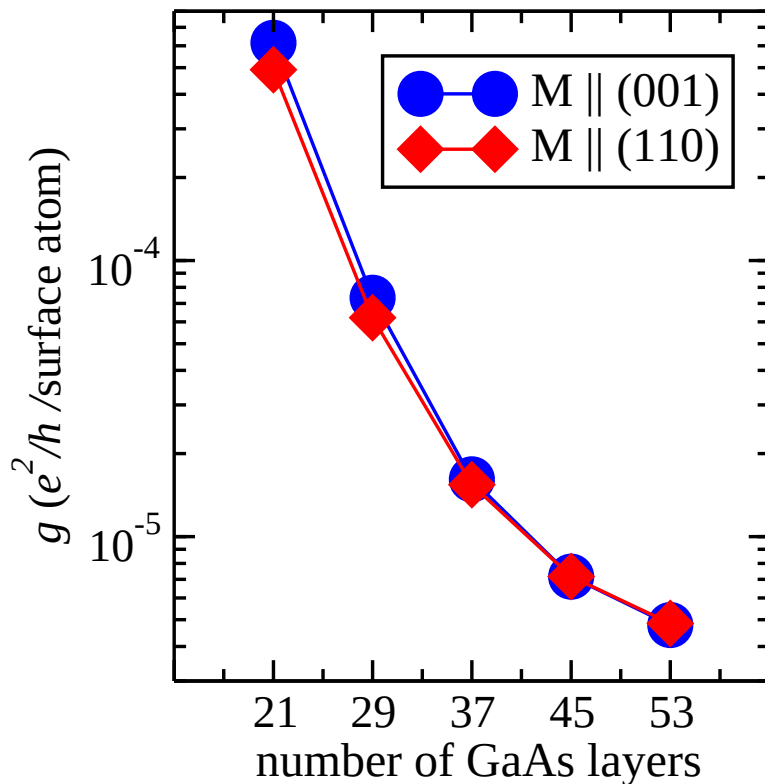


● difference of G largest for AP configuration

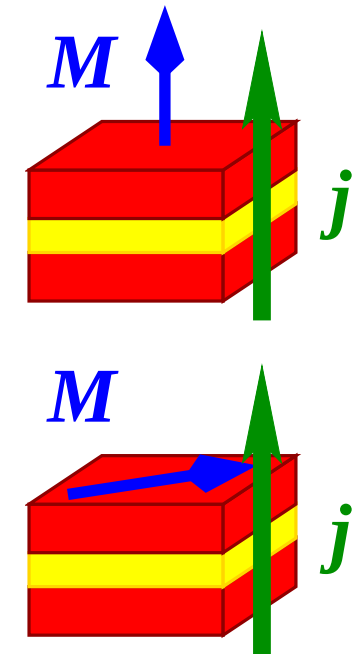
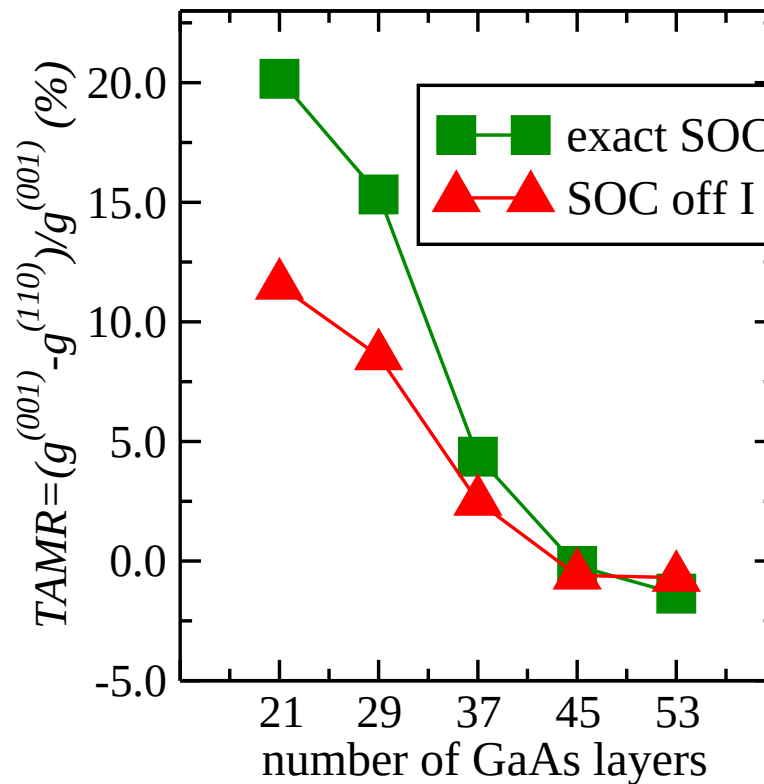
Fe/GaAs/Fe – TAMR calculations

Fe/*n*(GaAs)/Fe tunnelling junction – Ga-termination

● conductance

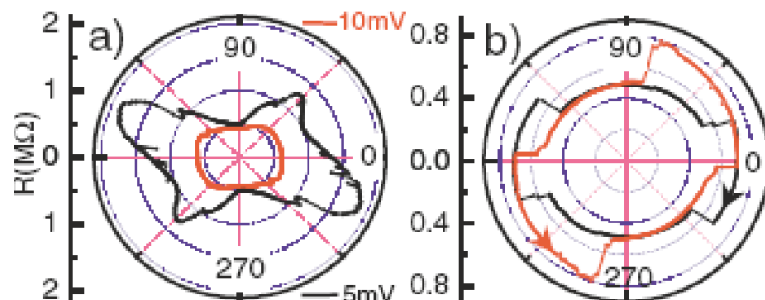
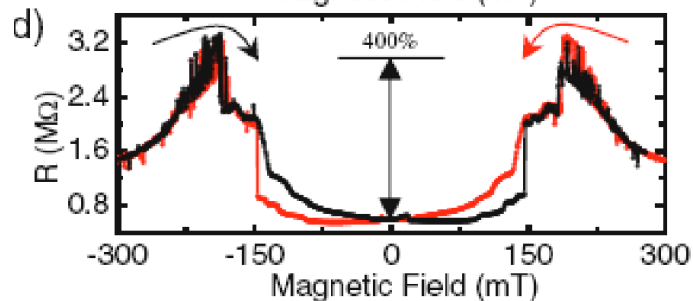
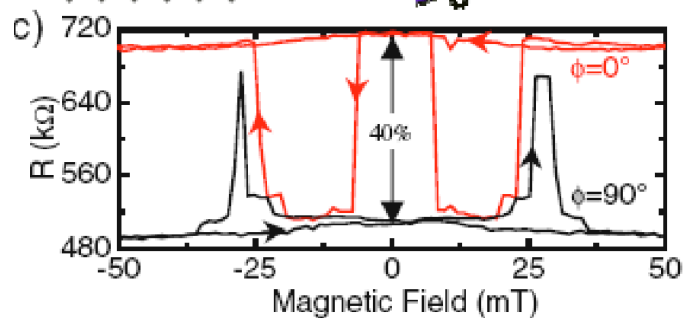
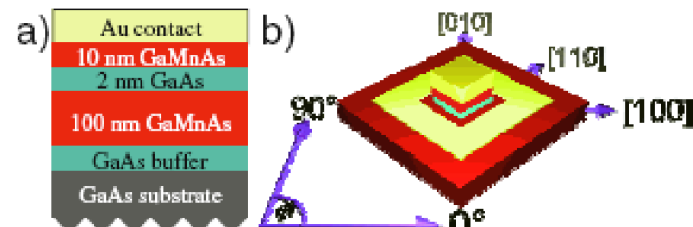


● TAMR



- strong TAMR effect for thin spacer
- suppressing SOC at the interface reduces the TAMR by $\approx 50\%$
- TAMR $\approx 6 \times$ larger than for As termination

Tunneling Anisotropic Magneto-Resistance

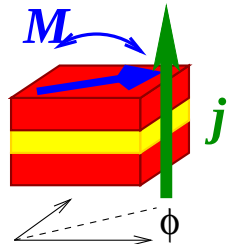
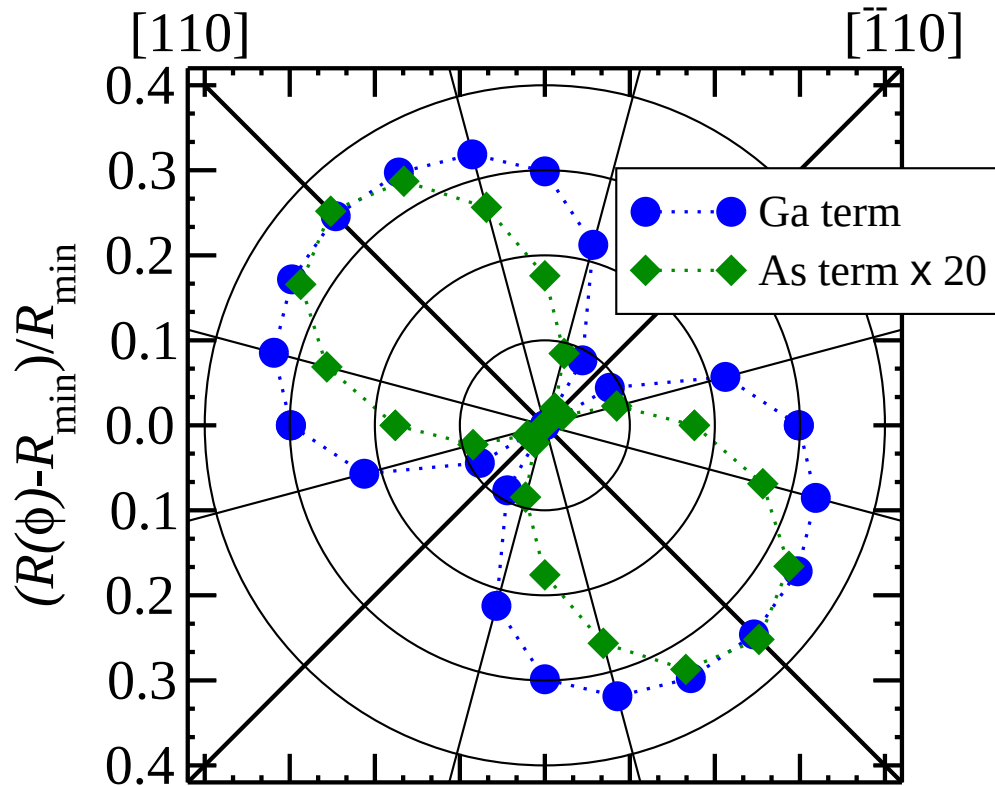


Gould *et al.* 2004, PRL 93, 117203
 Rüster *et al.* 2005, PRL 94, 027203

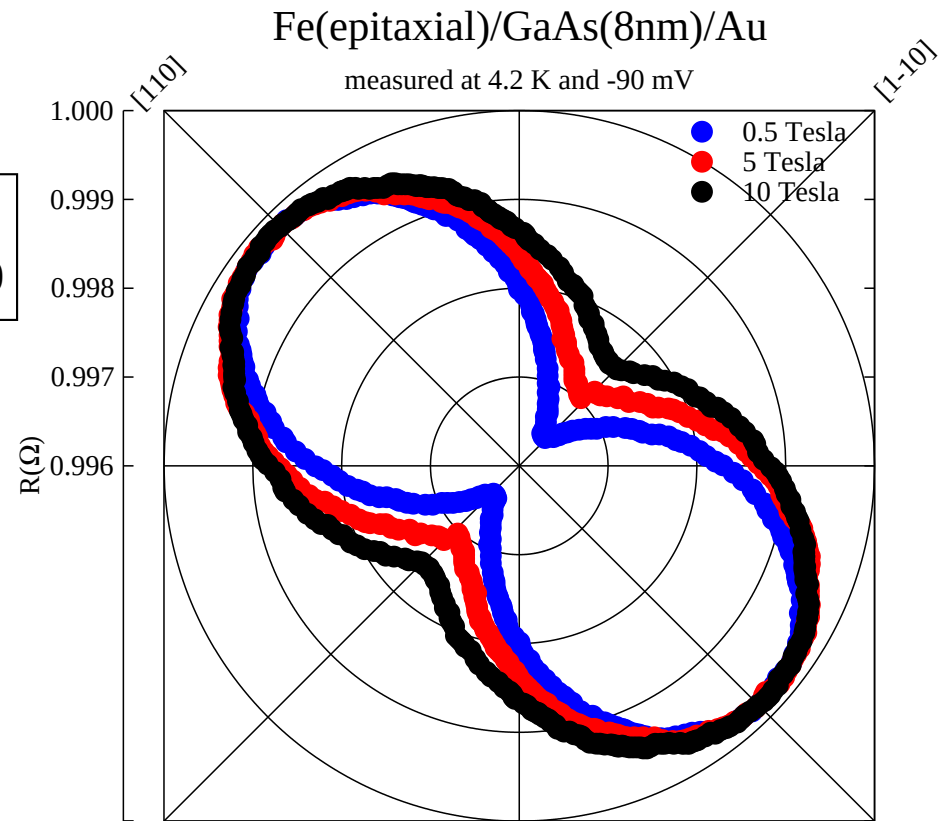
- **spin-valve** behaviour with only one (or two coupled) magnetic layer
- depending on ϕ width and sign change (**unlike** TMR)
- at saturation, the sample is a sensor of the **absolute direction** of $\vec{H}$
- modelling conductance g
 - $g \propto$ 'tunnelling' DOS
 - $\Delta \text{DOS} \equiv \text{DOS}(M \parallel x) - \text{DOS}(M \parallel y)$
- **spin-orbit coupling (SOC)** induced anisotropy

Fe/29(GaAs)/Au – in-plane TAMR

 **theory**



 **experiment**



Moser *et al.*, Univ. of Regensburg

Phys. Rev. Lett. **99** 056601 (2007)

 qualitative (but not quantitative) agreement between theory and experiment

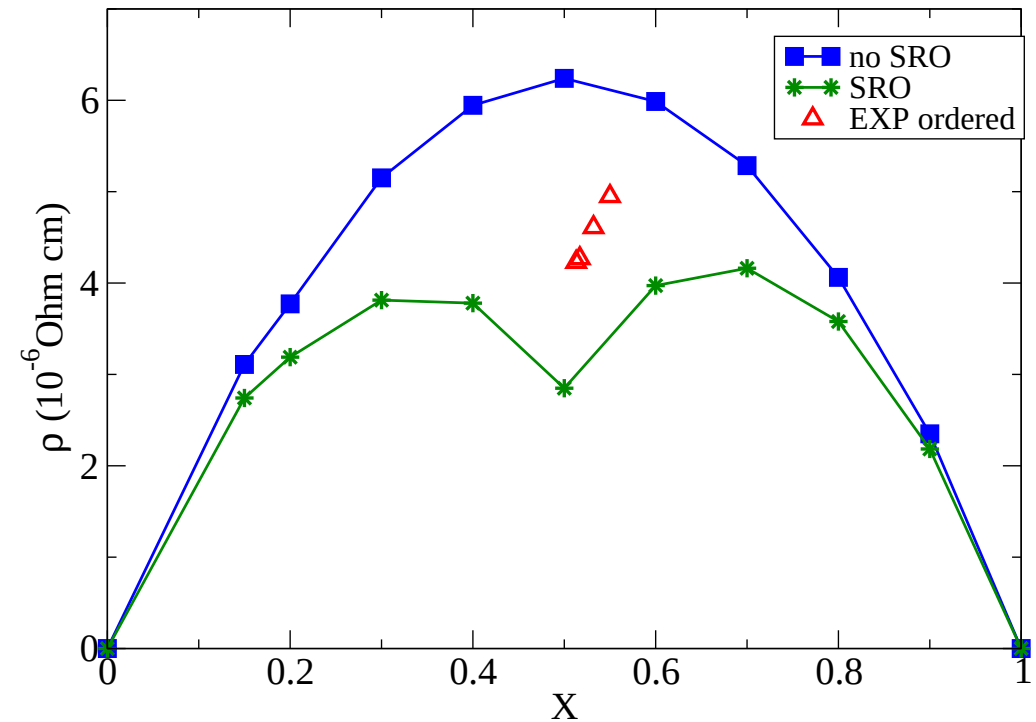
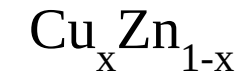
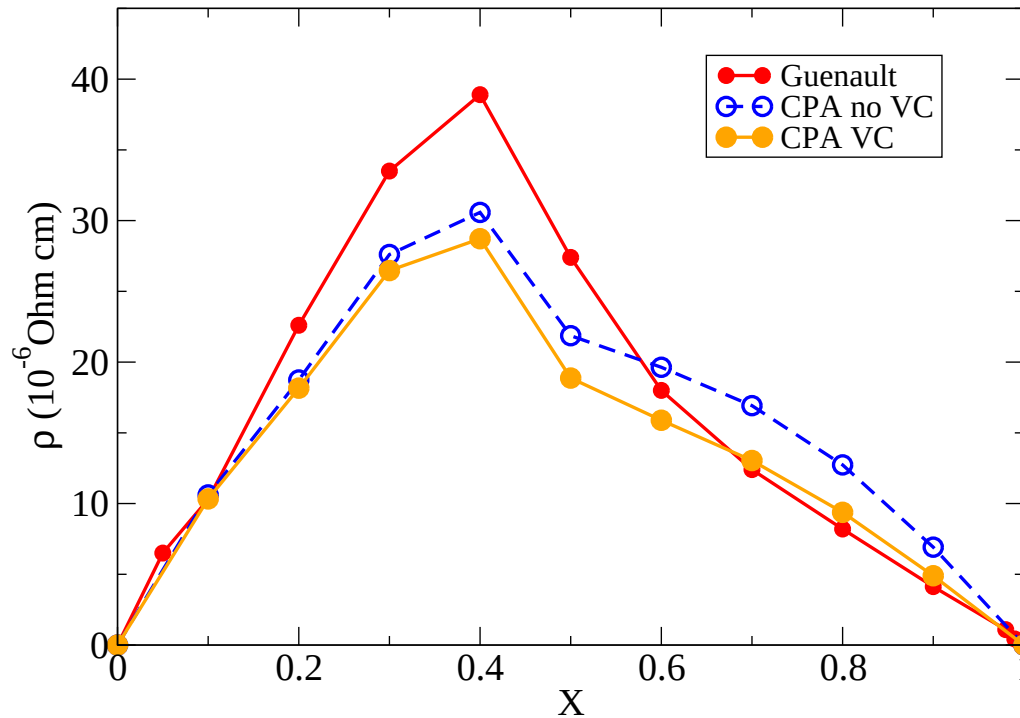
Středa, Smirčka, JPF (1975,1977)

$$\sigma_{\mu\nu} = -\frac{\hbar}{2\pi V} \int_{-\infty}^{+\infty} dE f(E) \text{Tr} \left\langle \hat{j}_{\mu} \frac{dG^{+}}{dE} \hat{j}_{\nu} (G^{+} - G^{-}) - \hat{j}_{\mu} (G^{+} - G^{-}) \hat{j}_{\nu} \frac{dG^{-}}{dE} \right\rangle$$

Crépieux, Bruno, PRB (2001), ($T = 0$)

$$\begin{aligned} \sigma_{\mu\nu} = & \frac{\hbar}{4\pi V} \text{Tr} \left\langle \hat{j}_{\mu} (G^{+} - G^{-}) \hat{j}_{\nu} G^{-} - \hat{j}_{\mu} G^{+} \hat{j}_{\nu} (G^{+} - G^{-}) \right\rangle \\ & + \frac{e}{4i\pi V} \text{Tr} \left\langle (G^{+} - G^{-}) (r_{\mu} \hat{j}_{\nu} - r_{\nu} \hat{j}_{\mu}) \right\rangle \end{aligned}$$

Residual resistivity of non-magnetic alloys

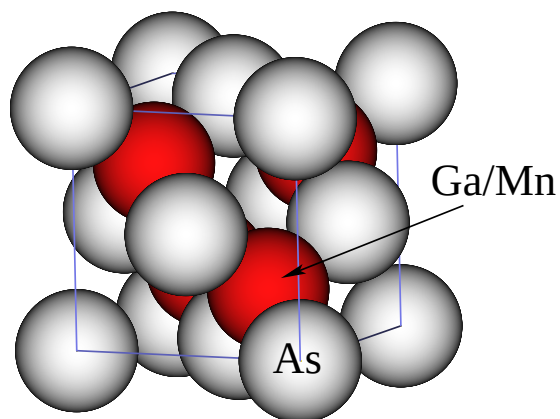
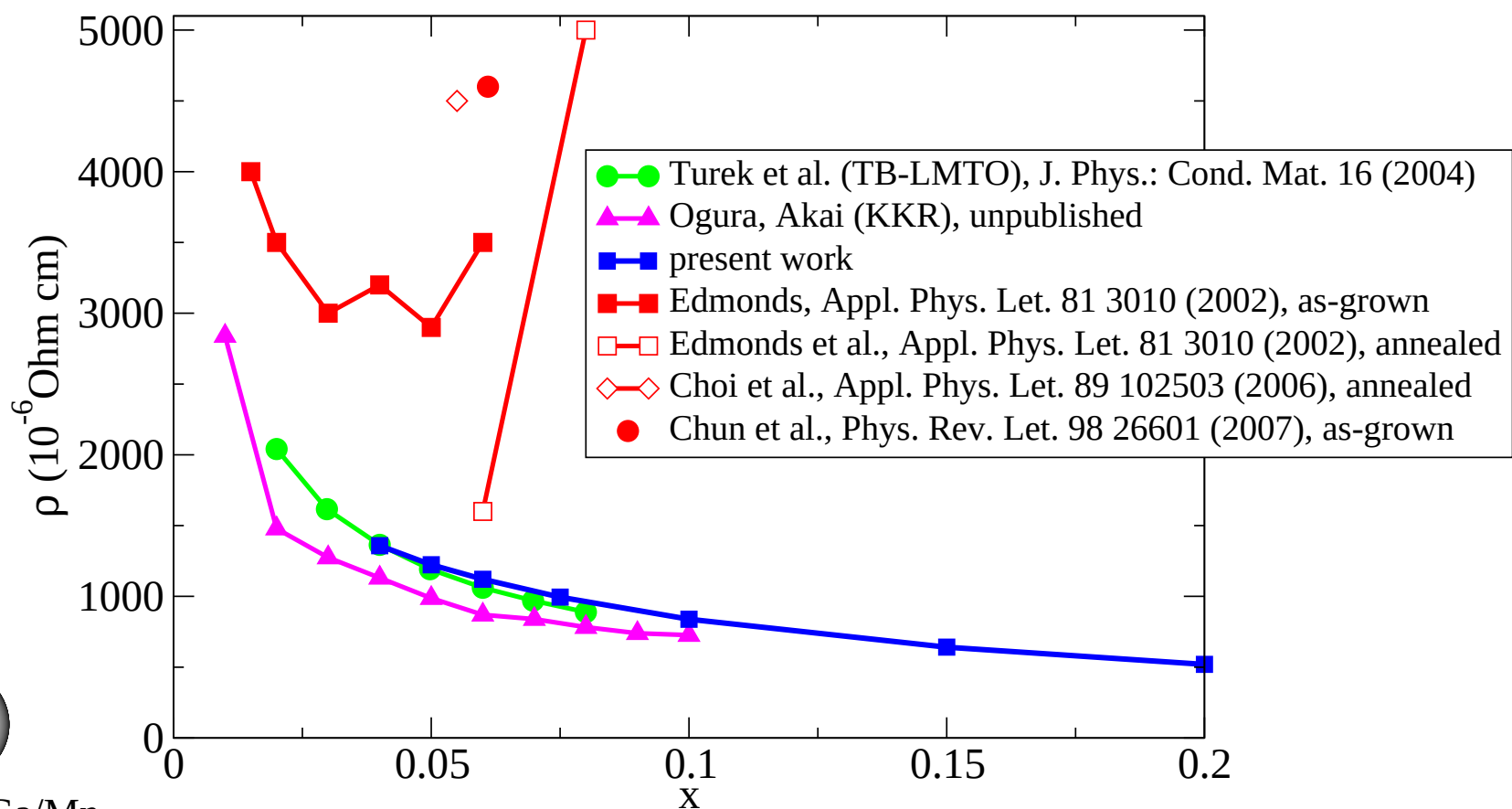


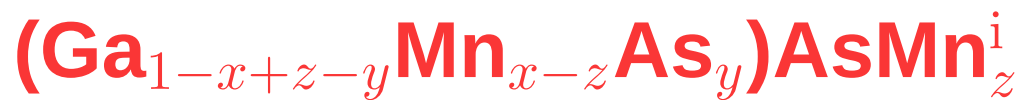
Experiment:

Guénault, Phil. Mag. 30, 641, 1974

W. Webb, Phys. Rev. 55, 297, 1938

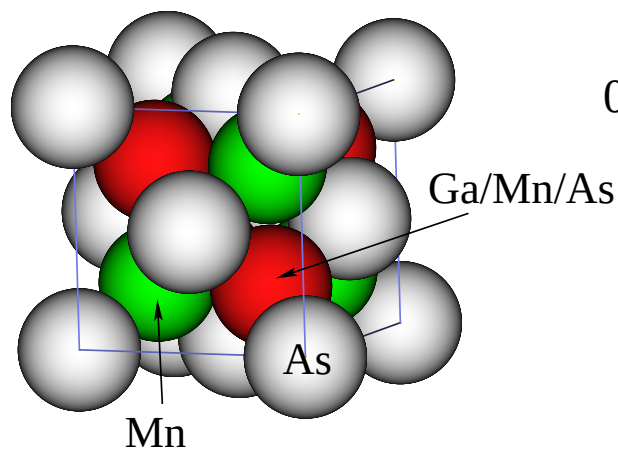
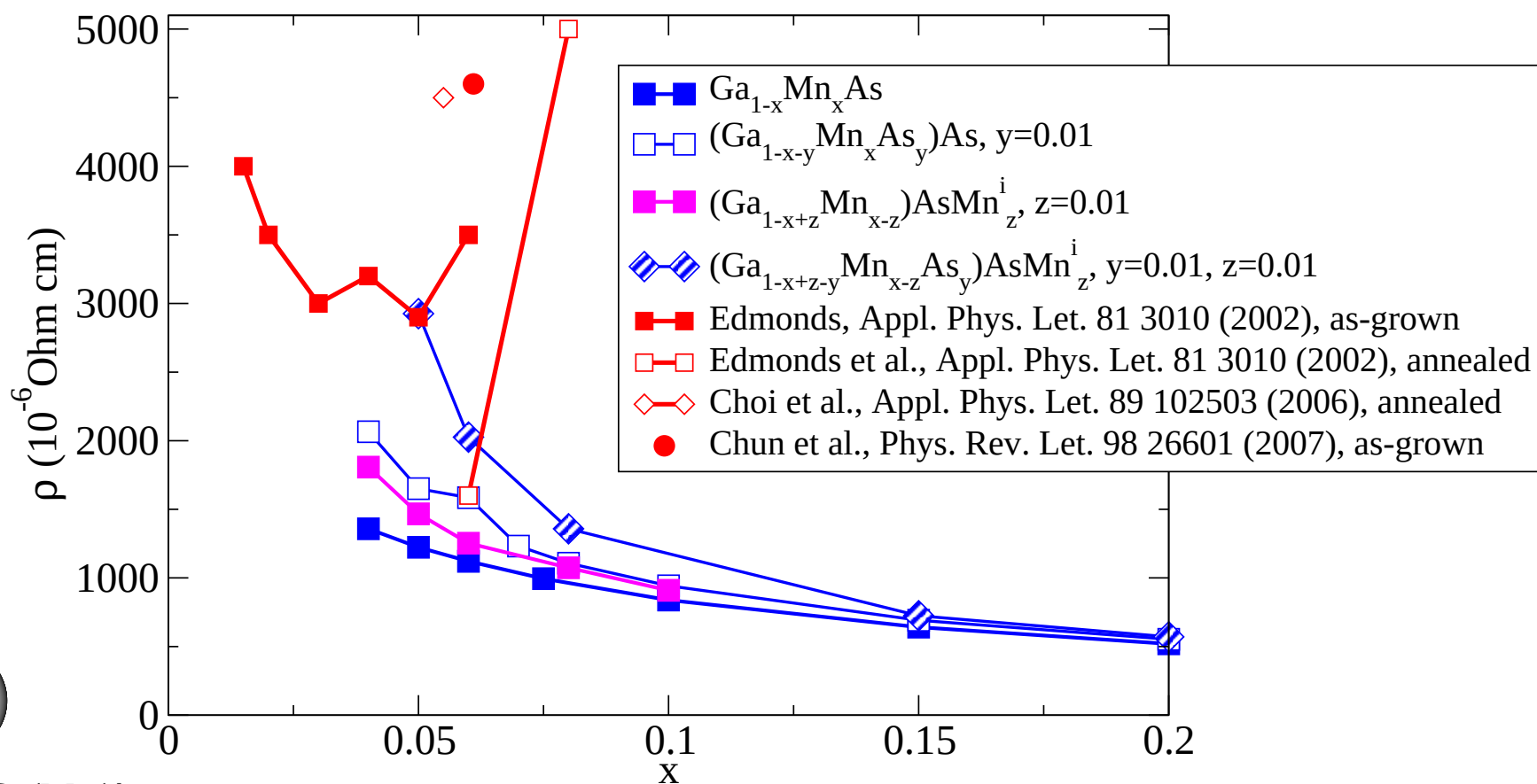
Residual resistivity comparison with experiment





Residual Resistivity

influence of antisite and interstitial occupation



Cubic ferromagnet $\vec{M} \parallel \vec{z}$

$$\rho = \begin{pmatrix} \rho_{xx} & \rho_{xy} & 0 \\ -\rho_{xy} & \rho_{xx} & 0 \\ 0 & 0 & \rho_{zz} \end{pmatrix}$$

Isotropic resistance $\bar{\rho} = \frac{1}{3}(2\rho_{\parallel} + \rho_{\perp})$

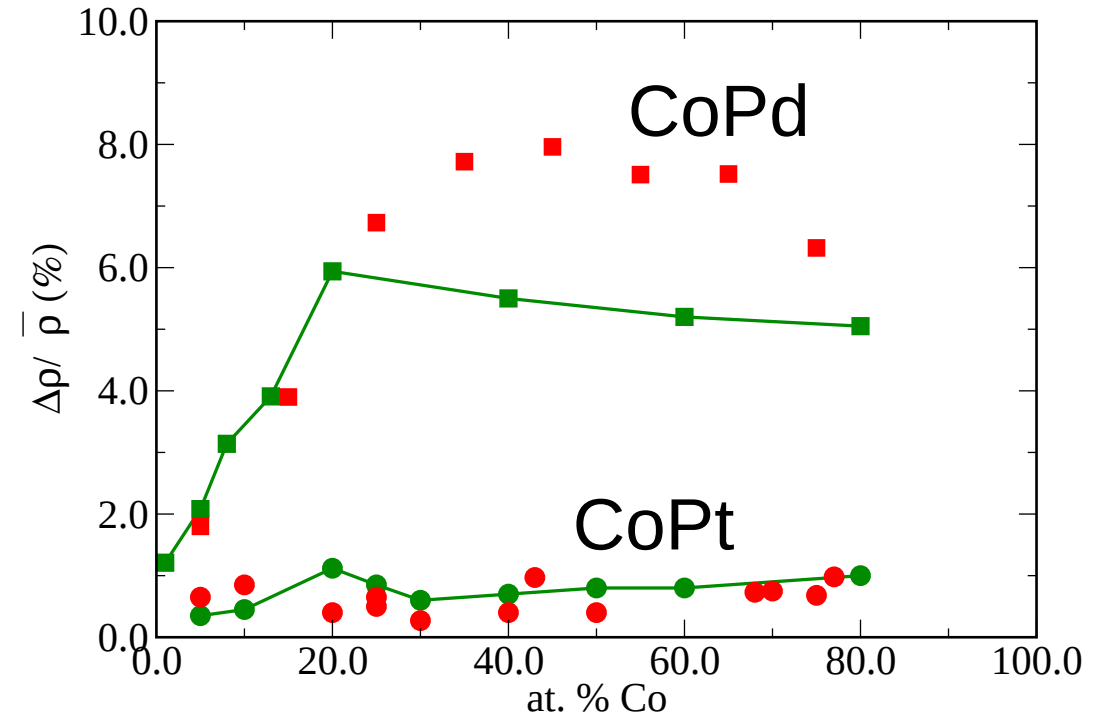
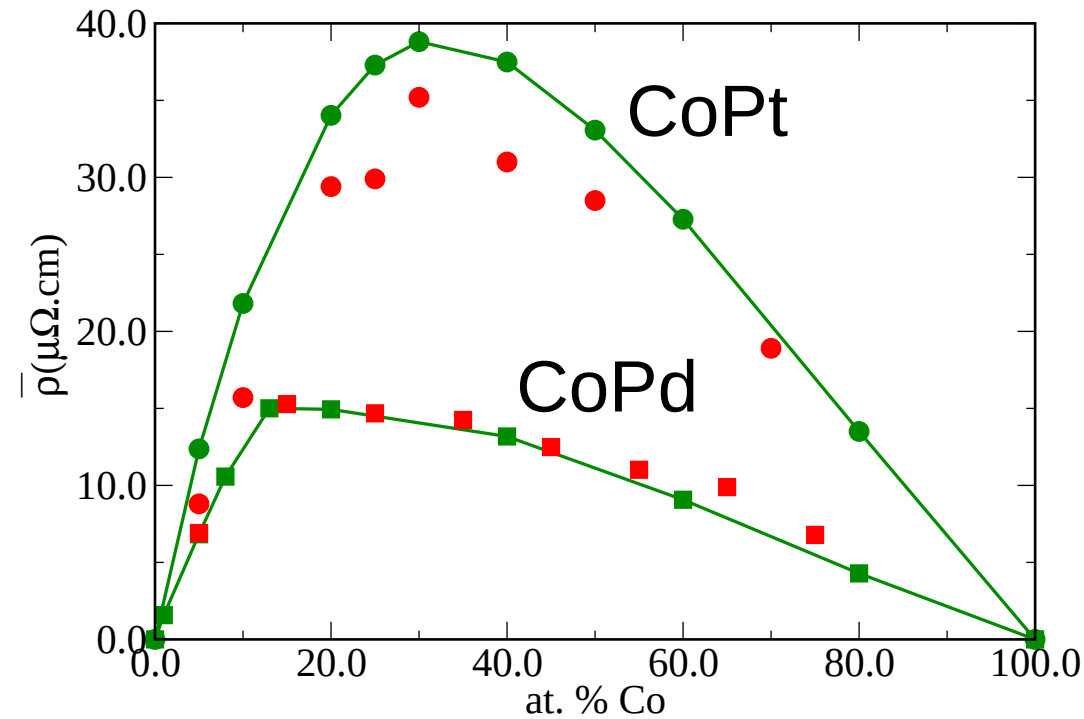
Anomalous magneto-resistance $AMR = \frac{\rho_{\parallel} - \rho_{\perp}}{\bar{\rho}}$

spontaneous Hall effect ρ_{xy}

Resistivity tensor of $\text{Co}_x\text{Pd}_{1-x}$ and $\text{Co}_x\text{Pt}_{1-x}$

Isotropic resistivity $\bar{\rho}$

Anomalous Magneto-Resistivity $\frac{\Delta\rho}{\bar{\rho}}$



Spin conductivity tensor

$$\vec{\Sigma}_{\mu\nu} = -\frac{\hbar}{2\pi V} \int_{-\infty}^{+\infty} dE f(E) \text{Tr} \left\langle \vec{J}_{\mu} \frac{dG^{+}}{dE} \hat{j}_{\nu} (G^{+} - G^{-}) - \vec{J}_{\mu} (G^{+} - G^{-}) \hat{j}_{\nu} \frac{dG^{-}}{dE} \right\rangle$$

$\vec{J}_{\mu}$ spin current operator for direction μ

j_{ν} electrical current operator for direction ν ($\vec{j} = ec\vec{\alpha}$)

Naive definition

$$J_{\mu}^z = \sigma_z j_{\mu}/e$$

Introduce auxiliary current that fulfills continuity equation

$$J_{\mu}^z = \frac{d}{dt} (\sigma_z r_{\mu}) = \sigma_z j_{\mu}/e + \left(\frac{d}{dt} \sigma_z \right) r_{\mu}$$

$\frac{d}{dt} \sigma_z$ spin-flip events due to SOC or non-collinear magnetism

Semi-relativistic: Niu (2005) and several others

Relativistic: Gyorffy, Szunyogh (2008)

Summary

- Fully relativistic description of transport within
 - Landauer-Büttiker
 - Kubo-Greenwood (and extensions)
- Role of SOC
 - influence on conductance G
 - in/out of plane anisotropy of TMR
 - in plane anisotropy of TMR-TAMR
 - source for bulk galvano-magnetotransport
- Formal basis for corresponding spin transport calculations