

Spin-polarized transport in ferromagnetic semiconductor / diffusive semiconductor / superconductor junctions

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Motivation

Superconductivity

Ferromagnet



Appearance of the new quantum phenomena

- Josephson current via ferromagnet
- Interplay between Andreev reflection or proximity effect and spin polarization

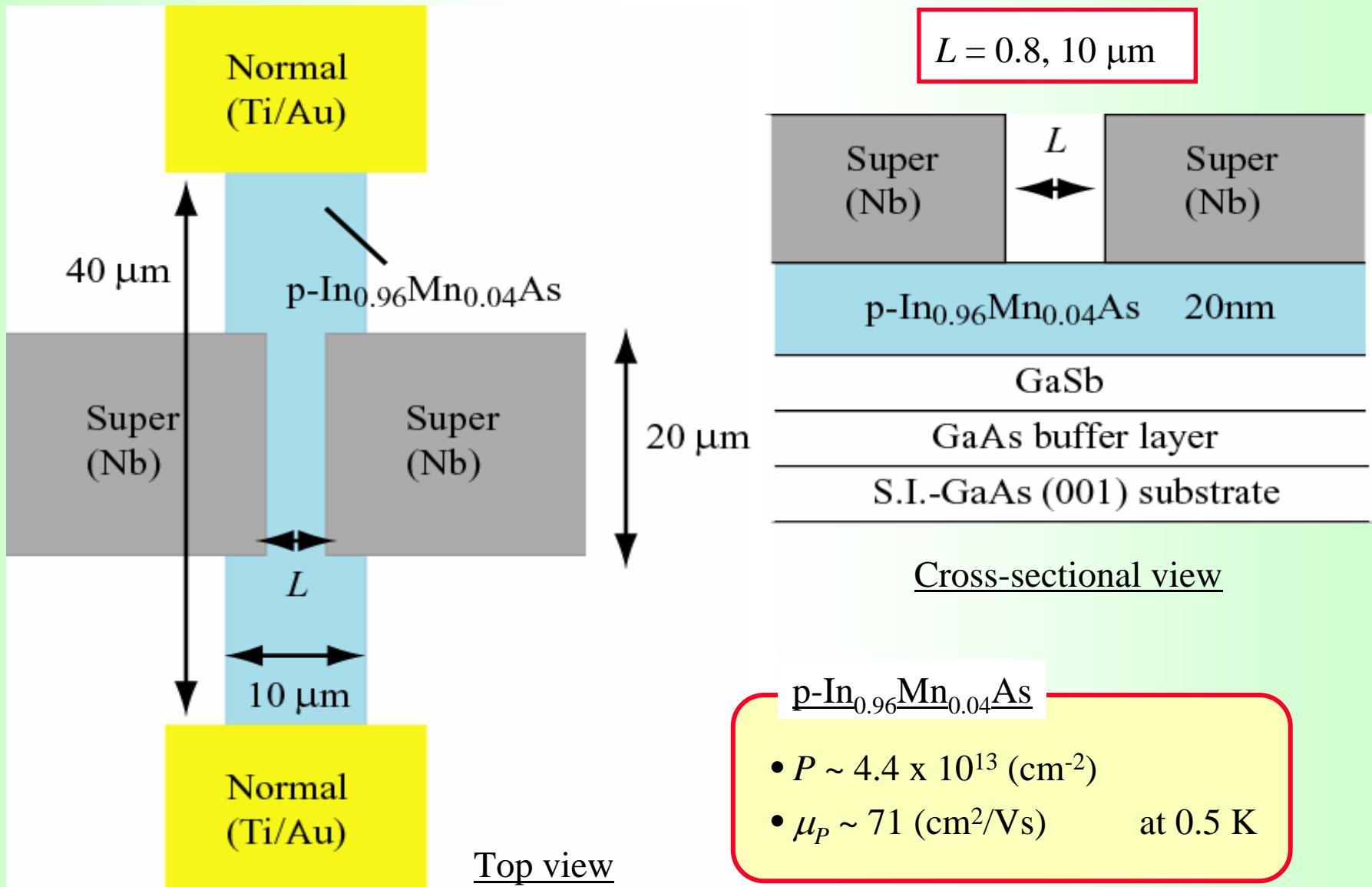
Nb/ferromagnetic p-InMnAs/Nb junction

S-F-S junction

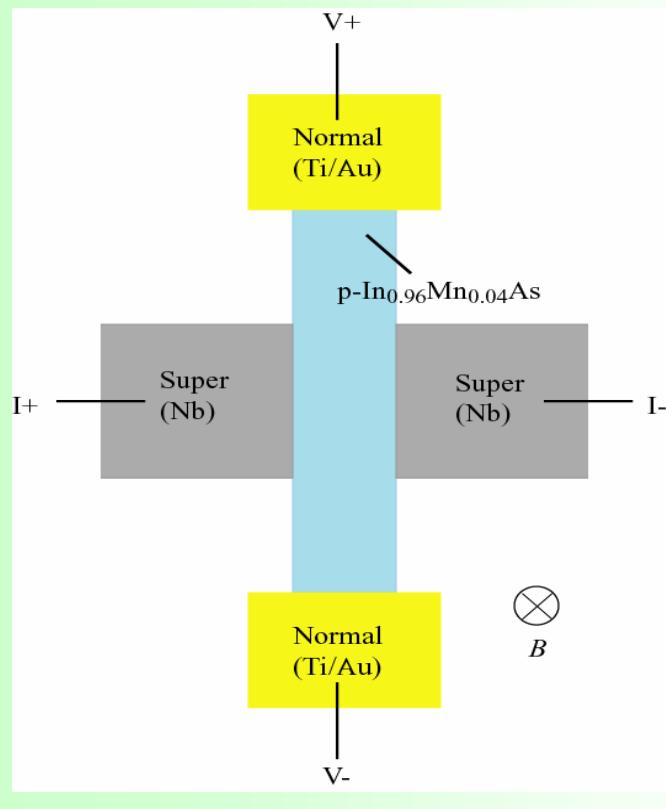
Nb/n-InAs/ferromagnetic p-InMnAs junction

S-N-F junction

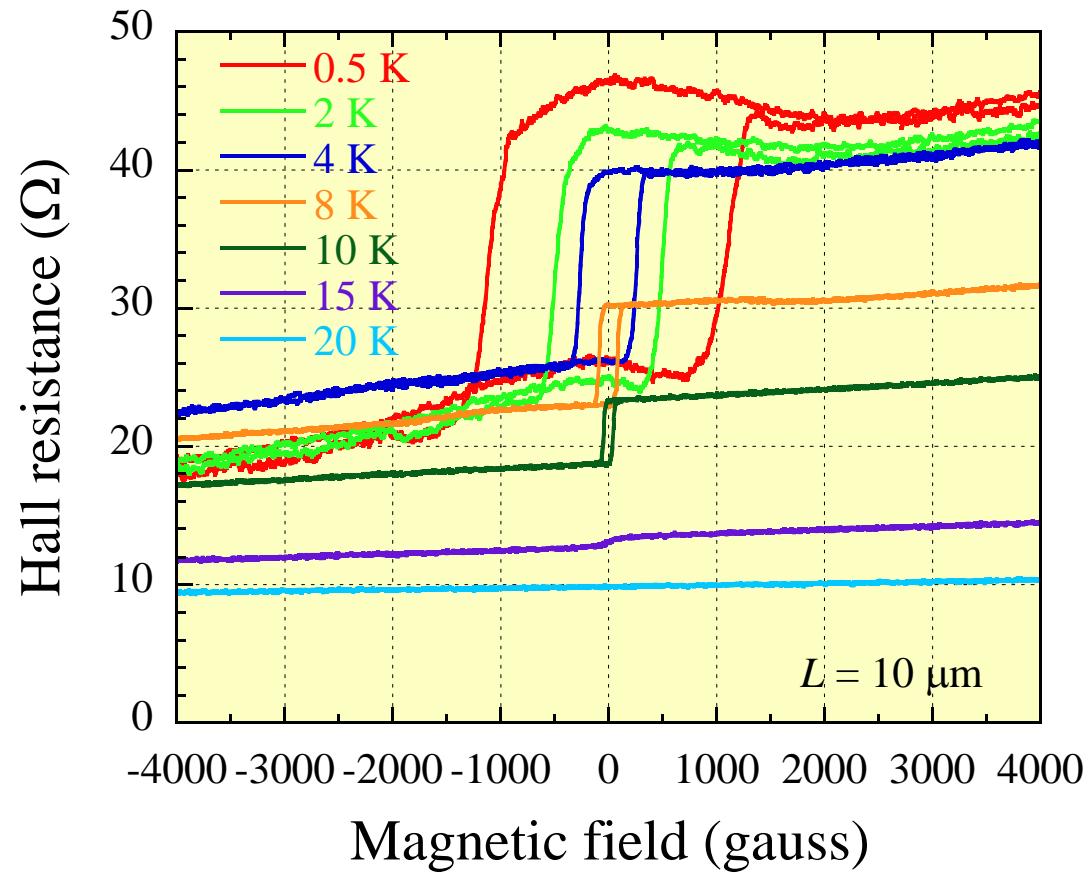
1. S-F-S junction Nb/p-InMnAs/Nb structure



Anomalous Hall effect

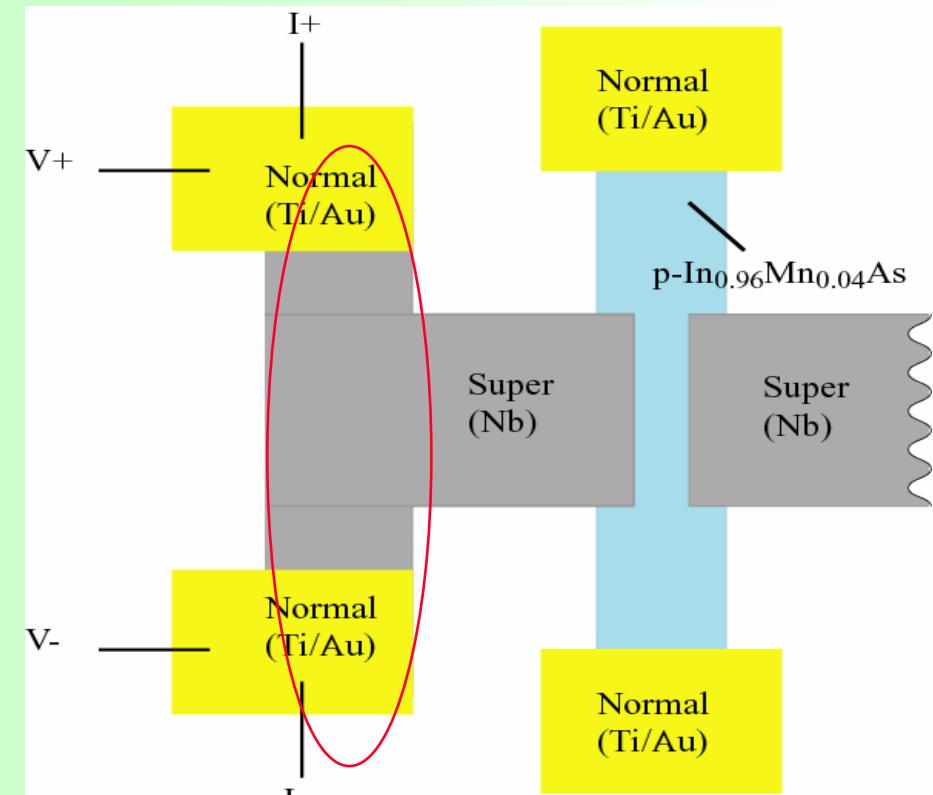


Measurement configuration

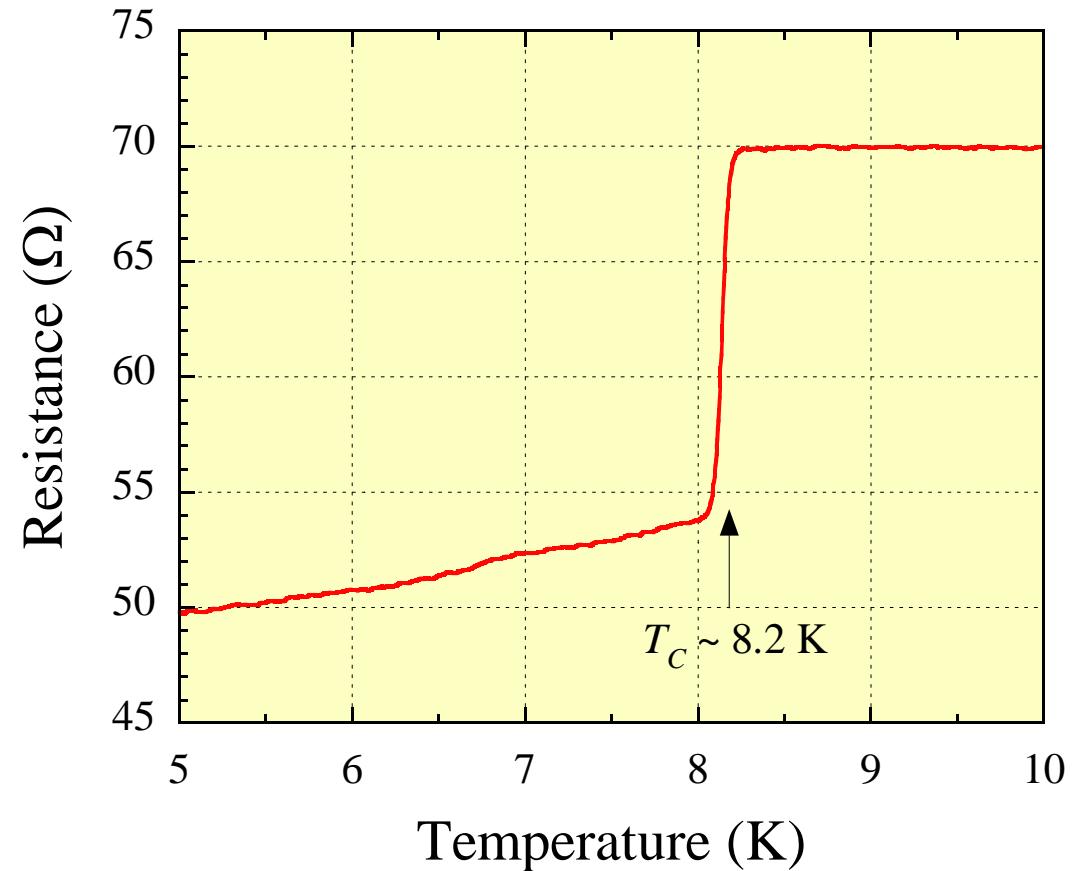


- Observation of anomalous Hall effect below ~ 15 K
- Reverse magnetic field is ~ 1000 gauss at 0.5 K.

T_C of Nb electrodes

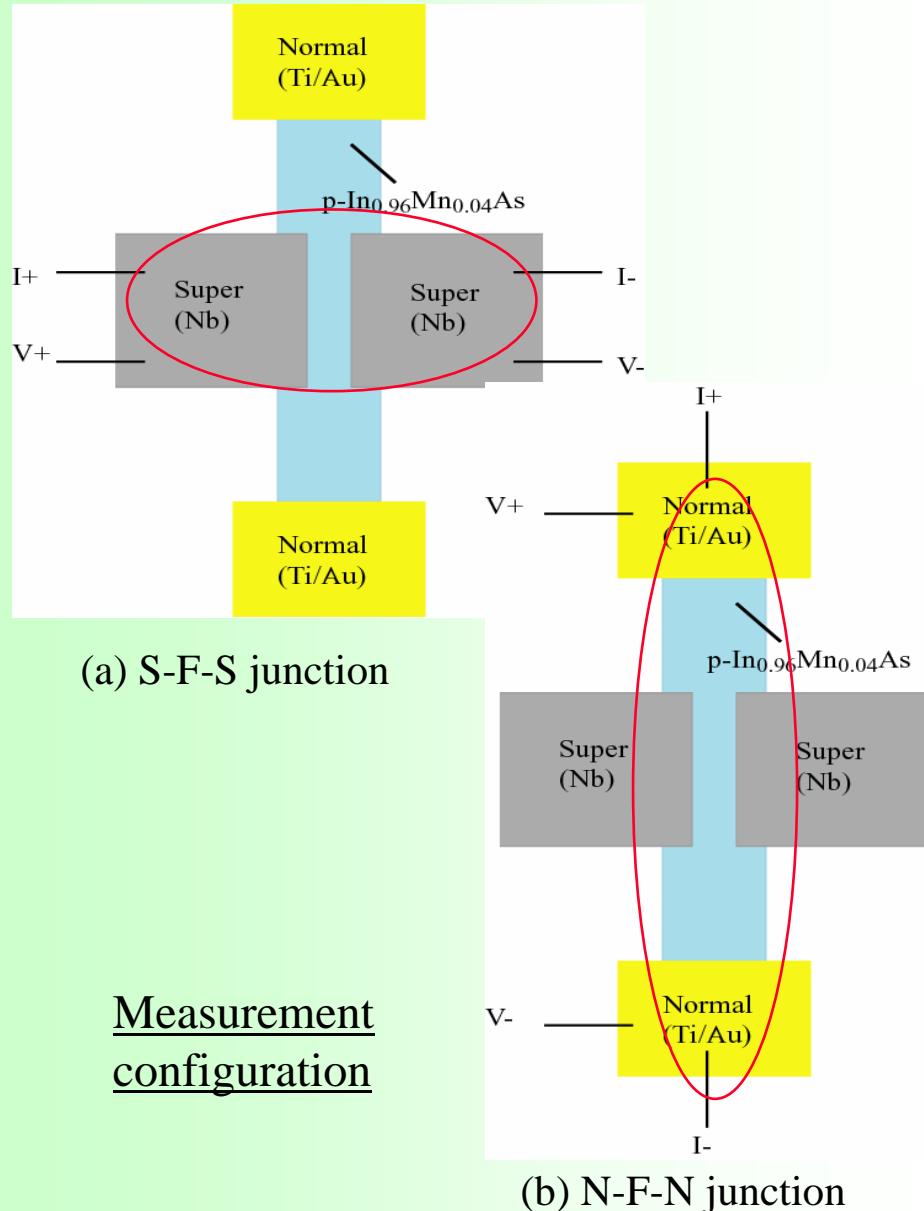


Measurement configuration

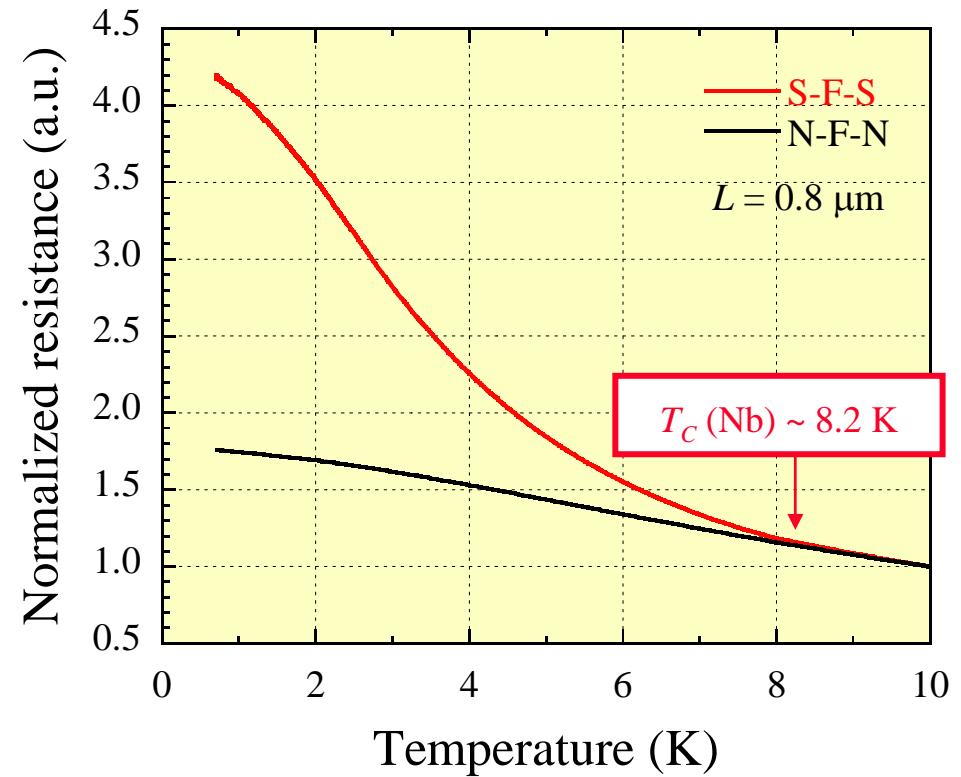


$T_C \sim 8.2$ K

Temperature dependence of resistance

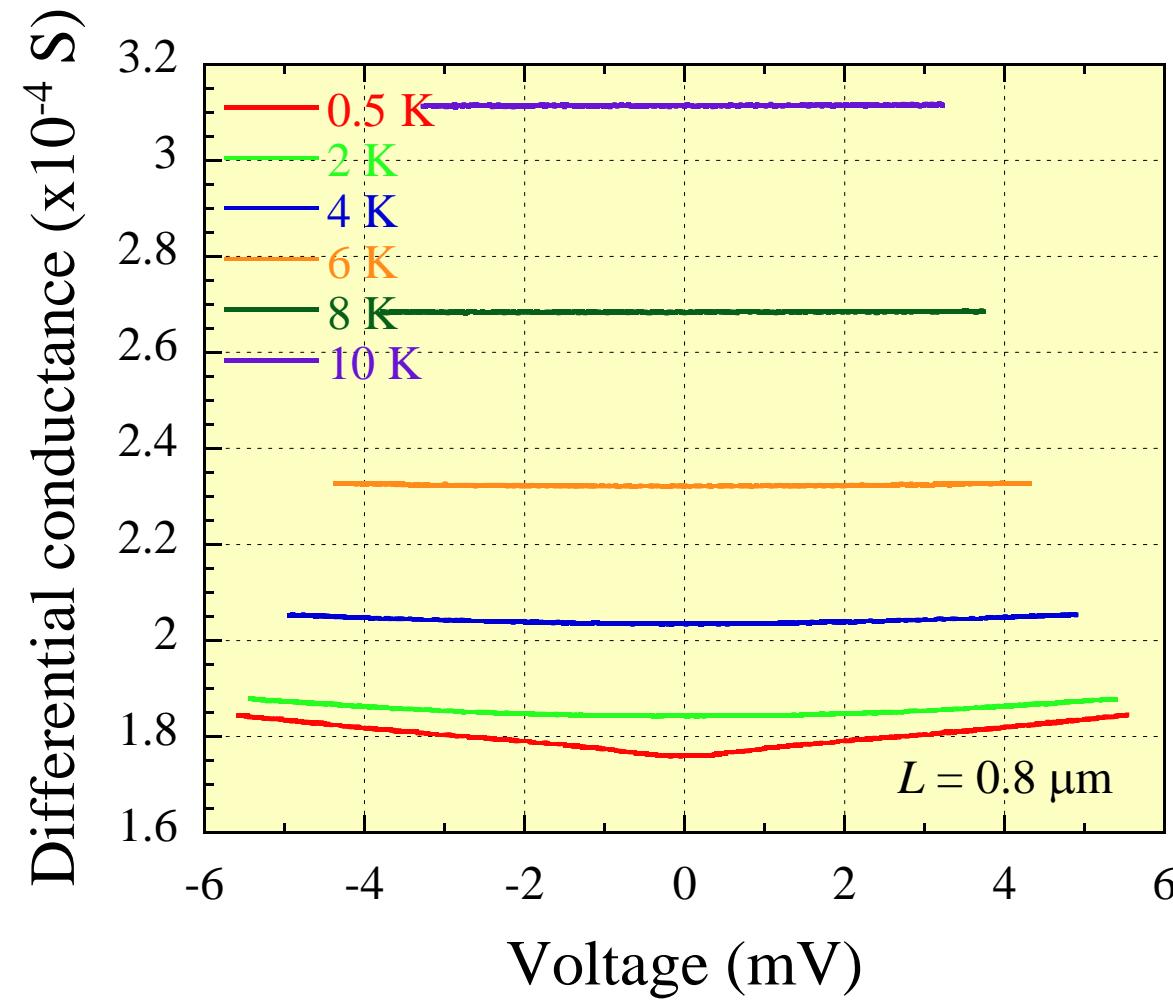


Measurement configuration



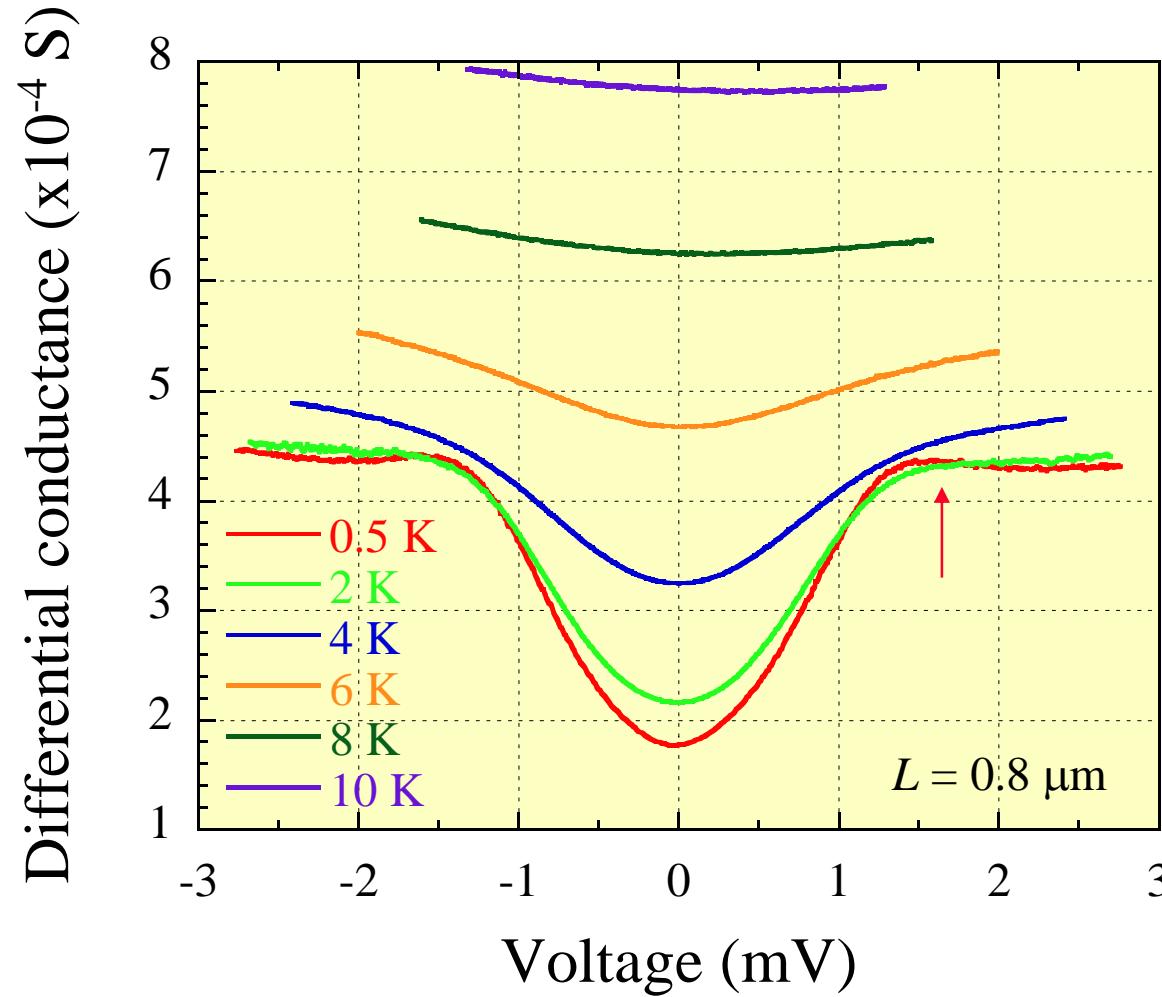
Below T_C of Nb, temperature dependence of resistance is completely different between the S-F-S and N-F-N junction.

Differential conductance in N-F-N junction



Although the weak tunneling behavior is observed in low temperatures, we have obtained nearly linear voltage-dependence.

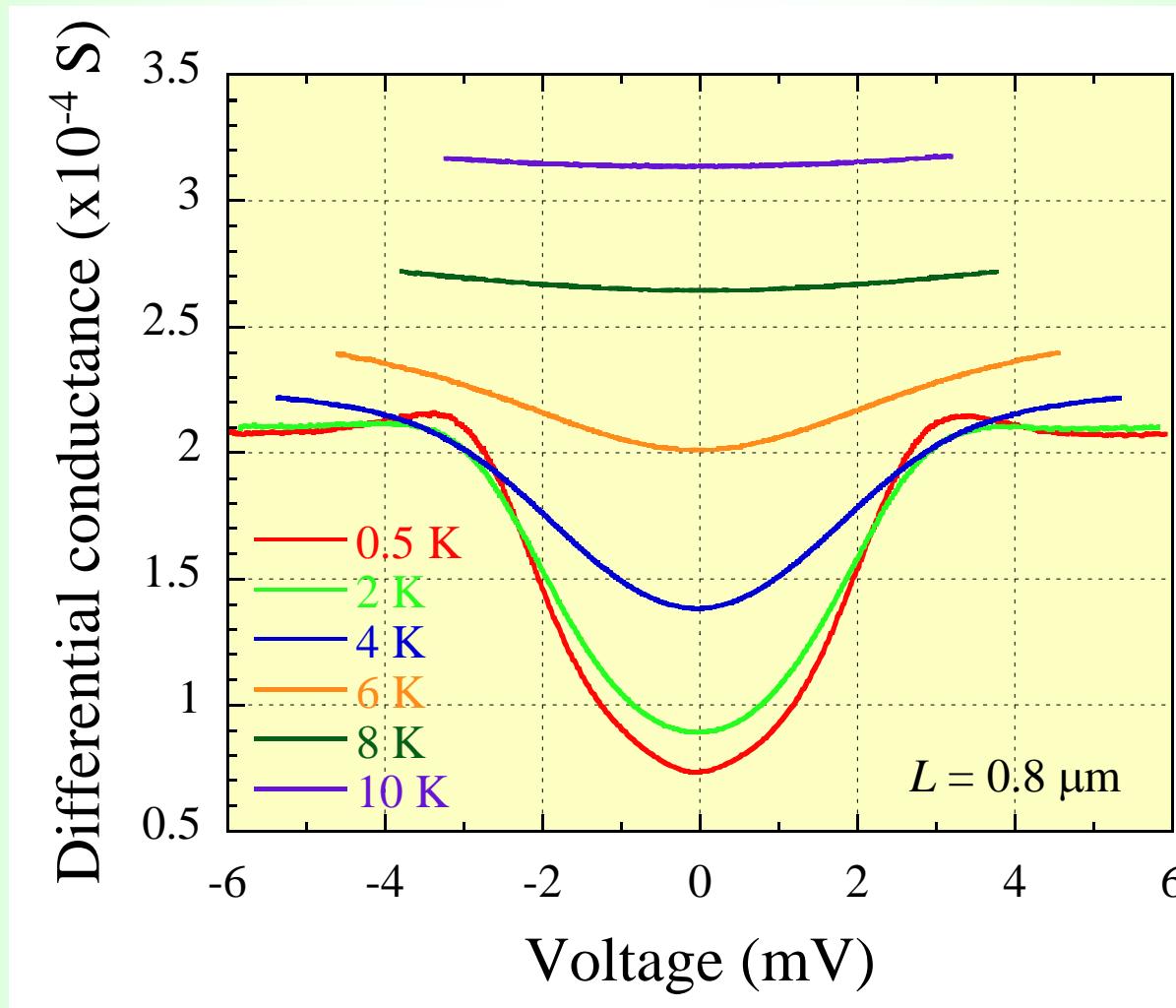
Differential conductance in S-F-N junction



We have obtained the conductance reduction within $V \sim 1.5$ mV.

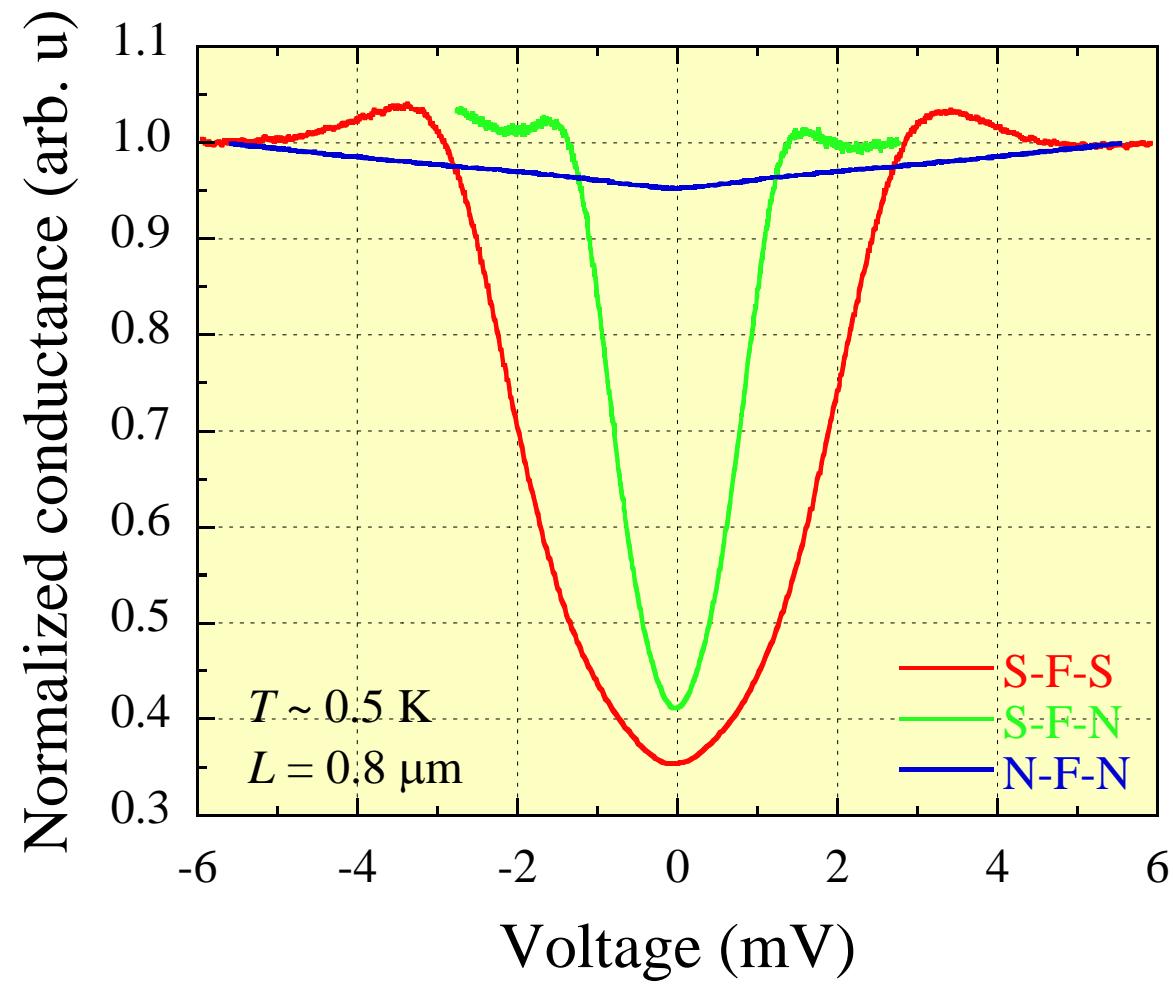
cf. Nb superconducting energy gap $\Delta_{\text{Nb}} \sim 1.5$ meV

Differential conductance in S-F-S junction



We have obtained the conductance reduction within $V \sim 3$ mV.

Comparison between all junctions



The conductance reduction is observed in **ONLY** S-F-S and S-F-N junctions.



The superconducting electrodes may affect conductance of the junction.

Discussion

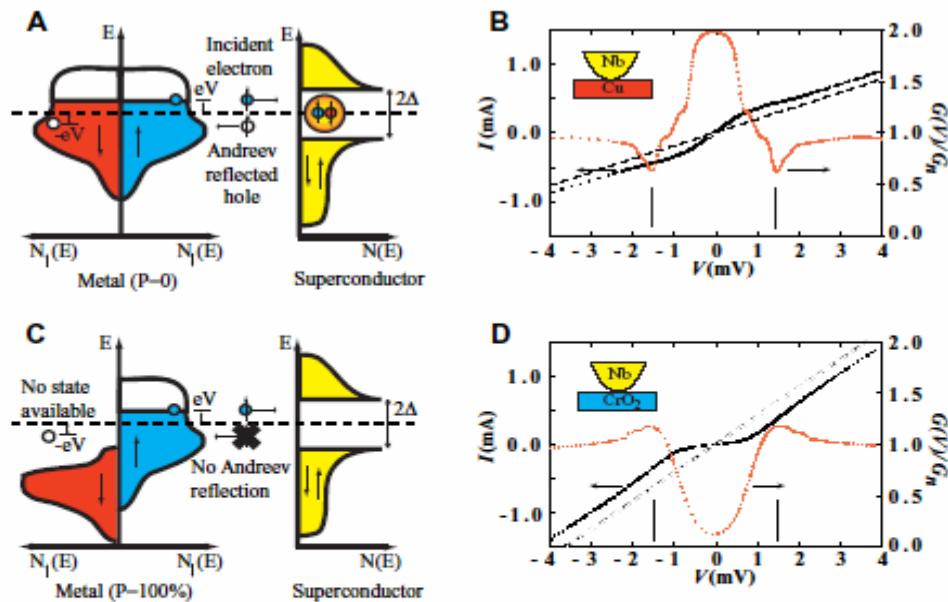


Fig. 1. Supercurrent conversion at the superconductor-metal interface for spin polarizations of $P = 0$ and $P \rightarrow 100\%$. (A) Schematic of the process for $P = 0$ when the Andreev reflection is unhindered by a spin minority population at E_F . The solid circles denote electrons and open circles denote holes. (B) Experimental measurement of the I - V and differential conductance dI/dV at $T = 1.6$ K via a superconducting Nb point contact on Cu. The vertical lines denote the bulk gap of Nb: $\Delta(T = 0) = 1.5$ meV. The dashed line is the normal state I - V for a conductance of $G_n = 0.194 \text{ ohm}^{-1}$. (C) Schematic of process for $P \rightarrow 100\%$ when there is no supercurrent conversion at the interface. (D) Experimental I - V and dI/dV at $T = 1.6$ K via the Nb point contact on CrO_2 . The dashed line is the normal state I - V for a conductance of $G_n = 0.417 \text{ ohm}^{-1}$.

R. J. Soulen Jr. et al., Science 282, p.86 (1998)

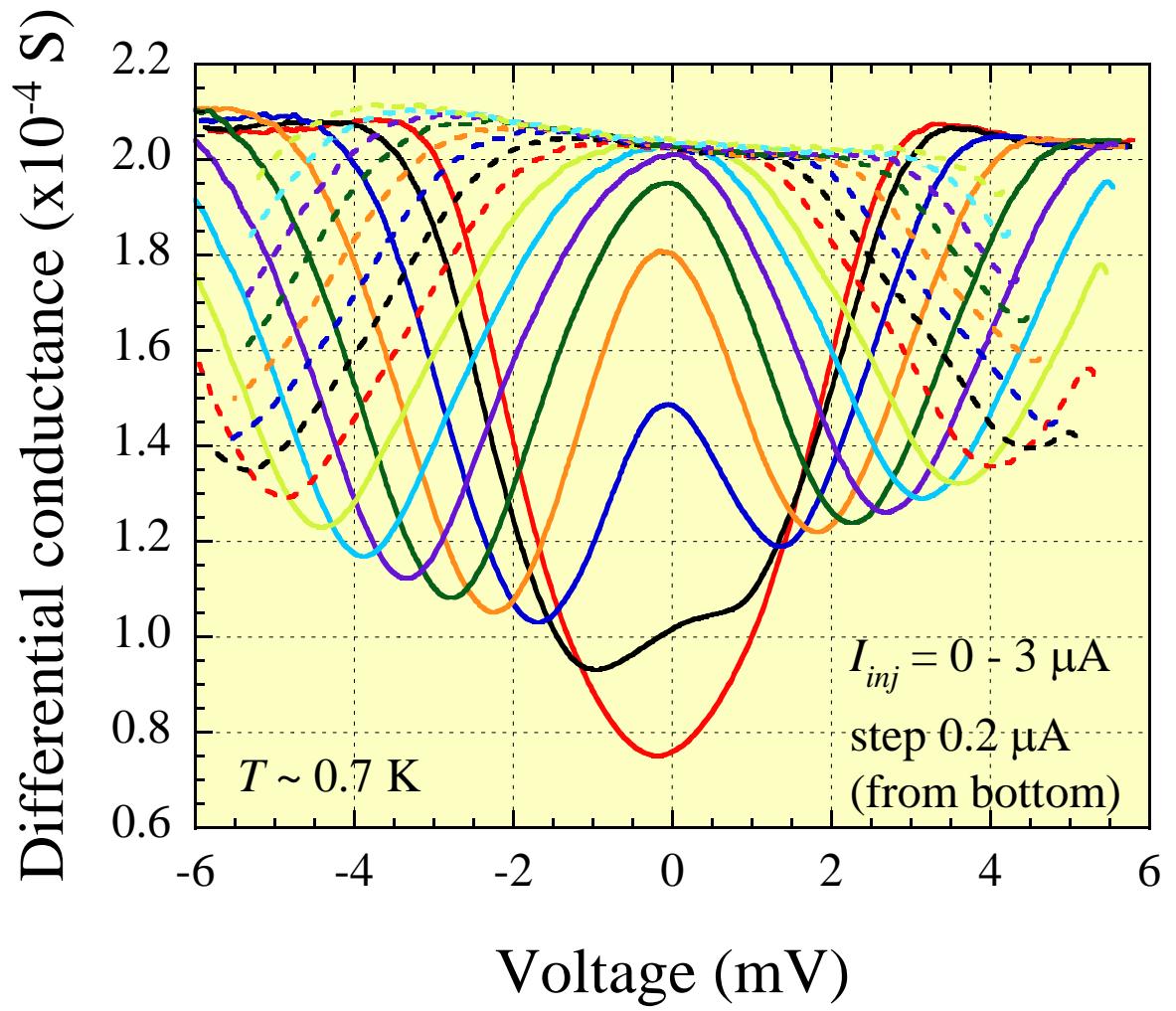
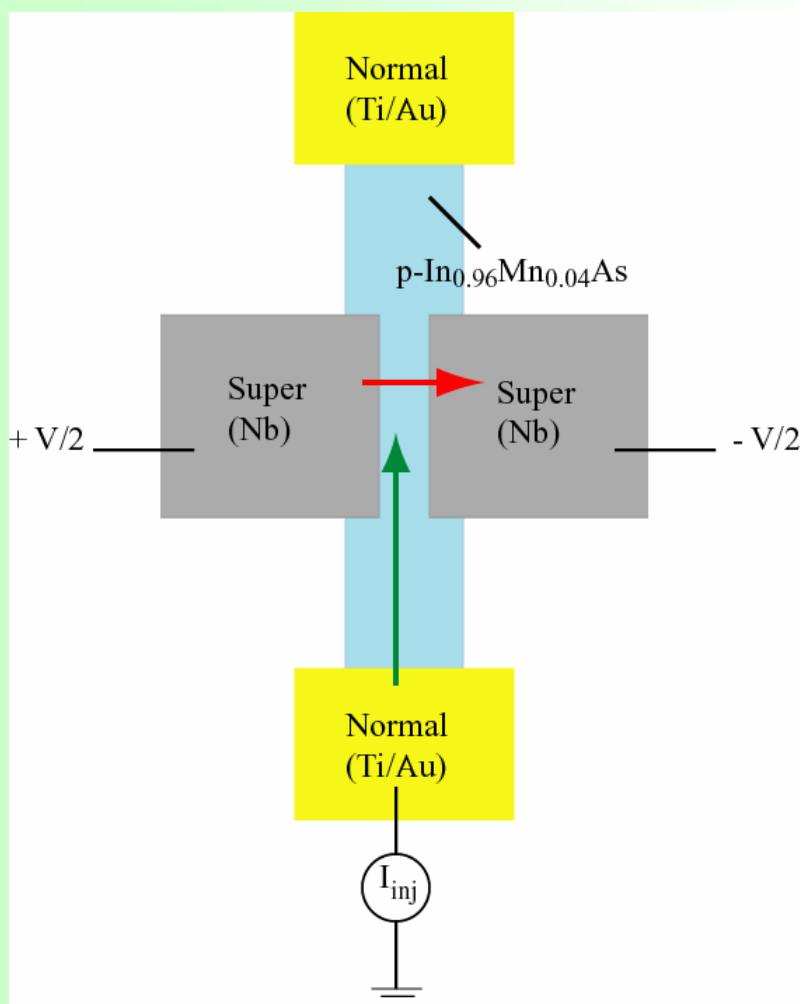
In Andreev reflection process, the incident electron requires the opposite spin electron to be removed from the N region for conversion to Cooper pair.



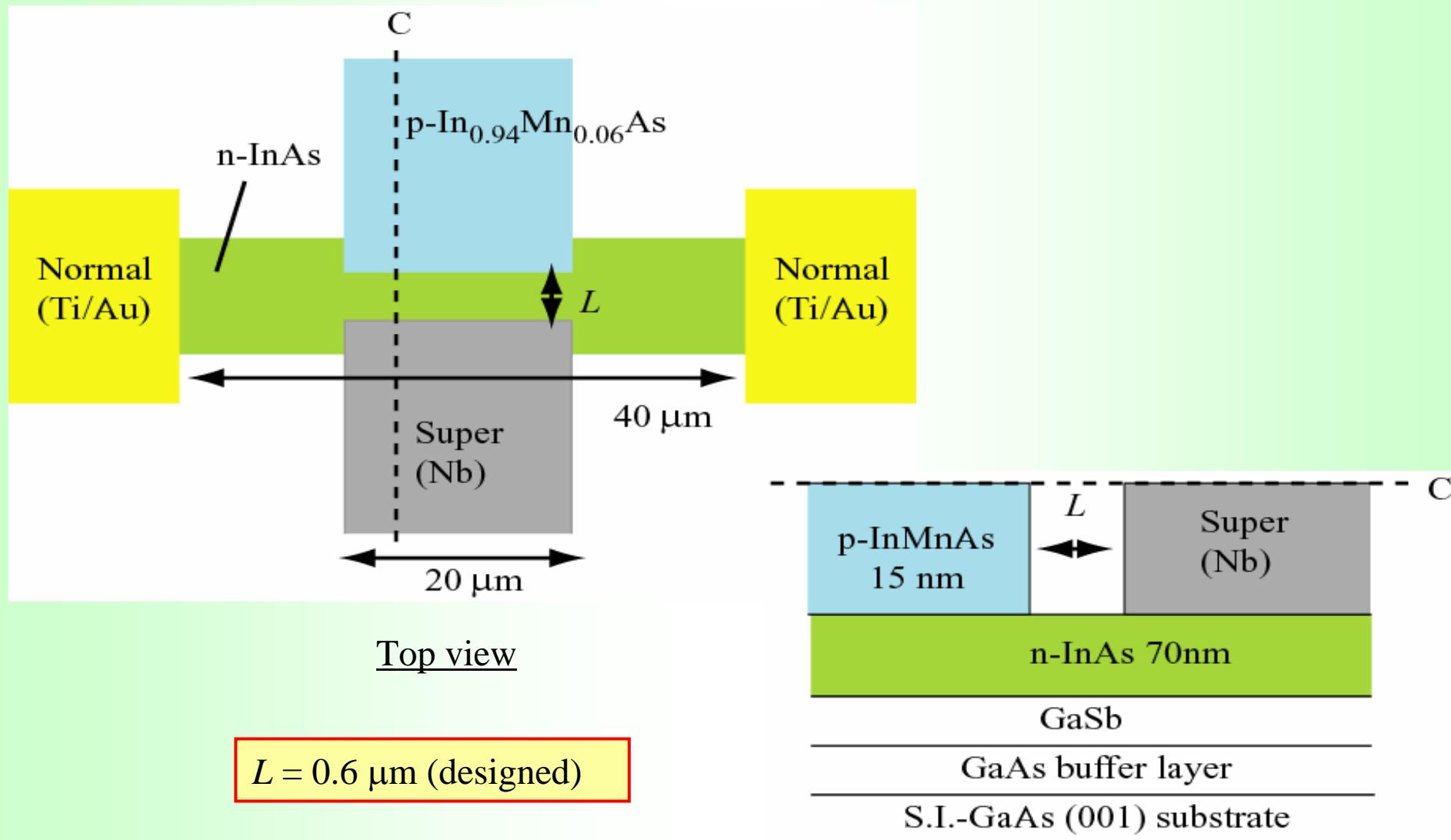
In case of S-F junctions, Andreev reflection is limited by the minority spin population.

Our experimental results can be qualitatively understood by considering the suppression of Andreev reflection due to spin polarization in p-In_{0.96}Mn_{0.04}As.

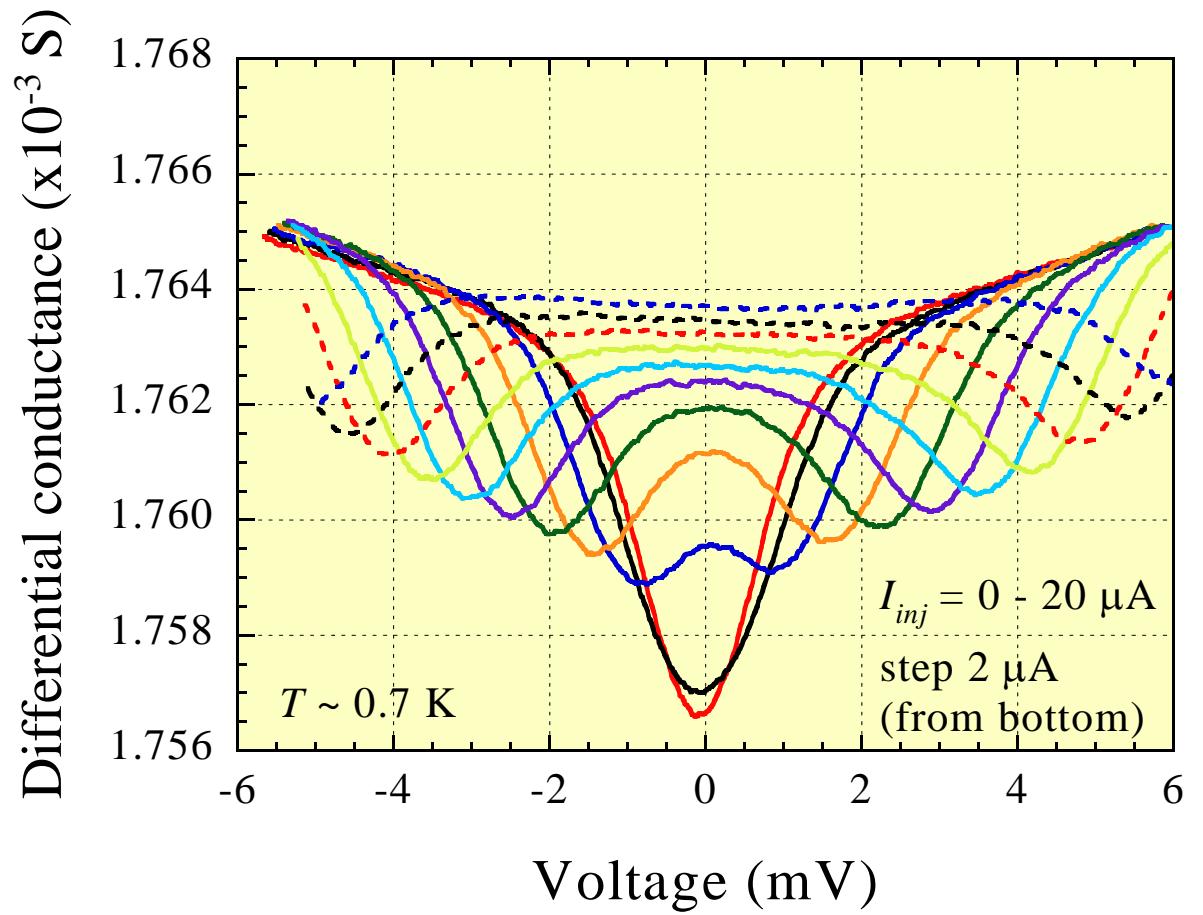
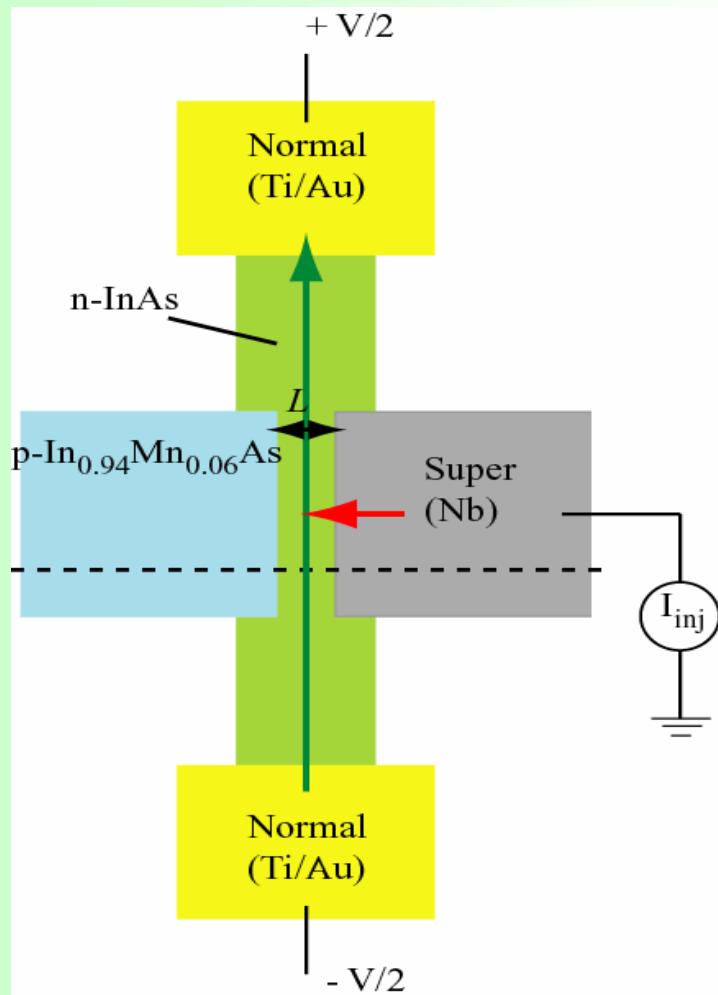
Current injection to SFS JJ



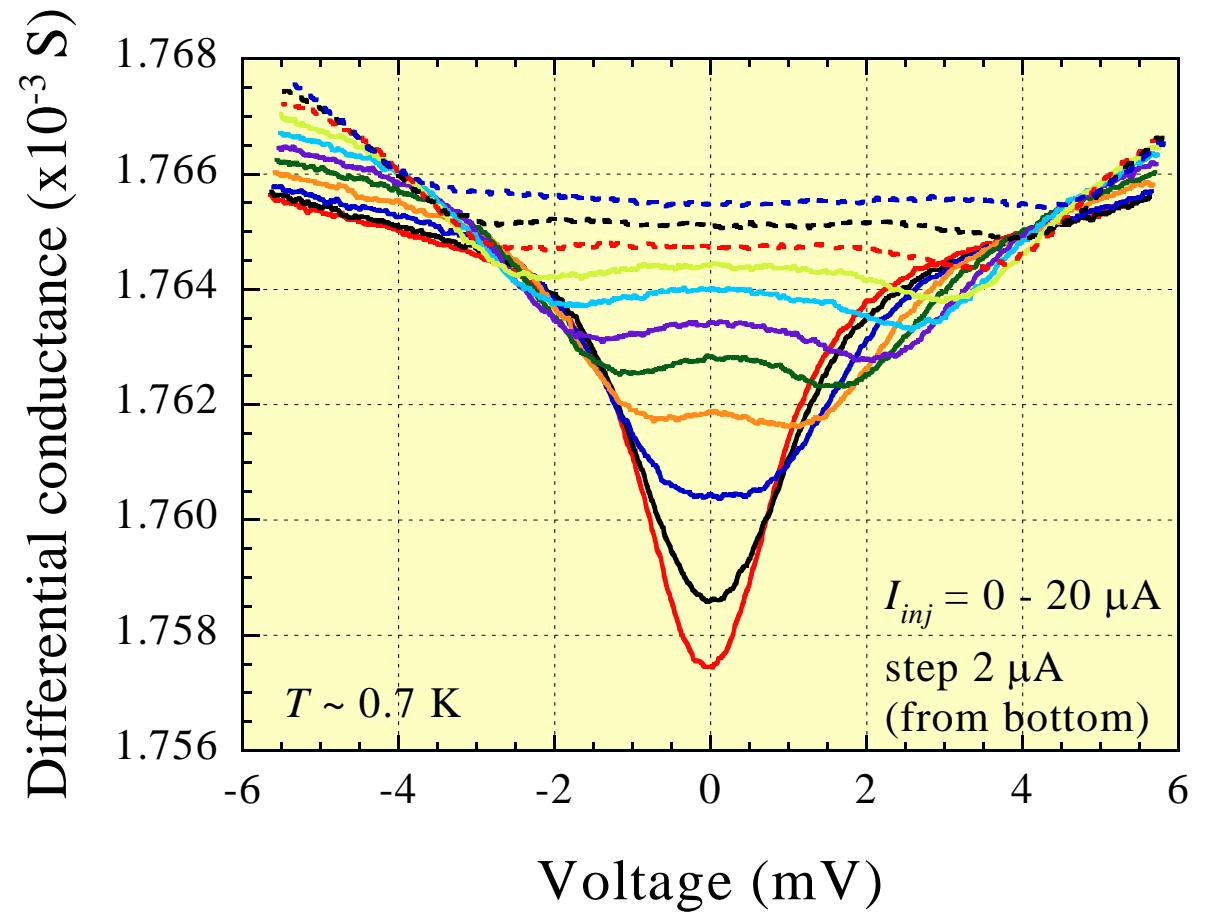
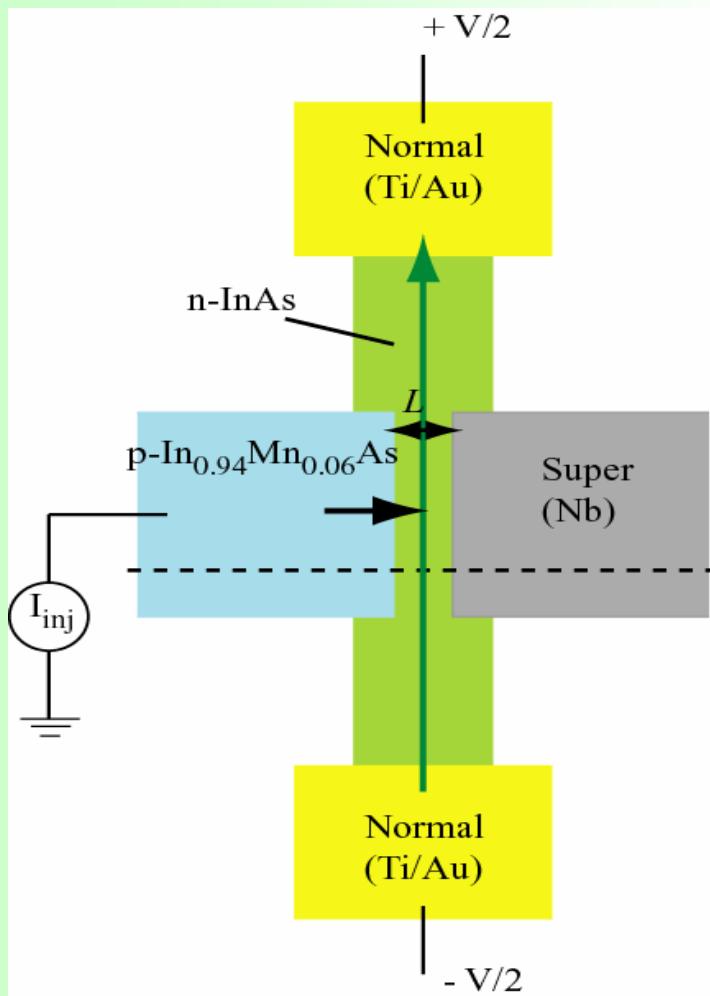
2. S-N-F junction p-InMnAs/n-InAs/Nb structure



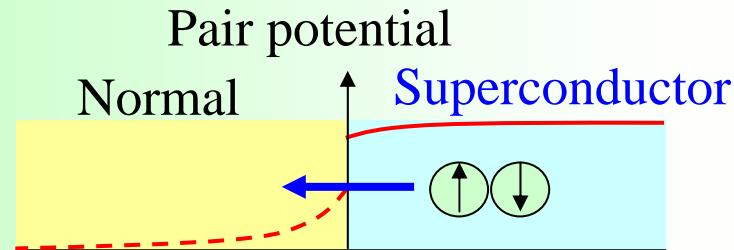
Current injection from Nb



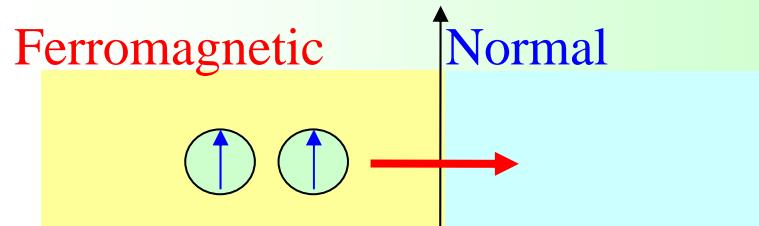
Current injection from p-InMnAs



Proximity Effect

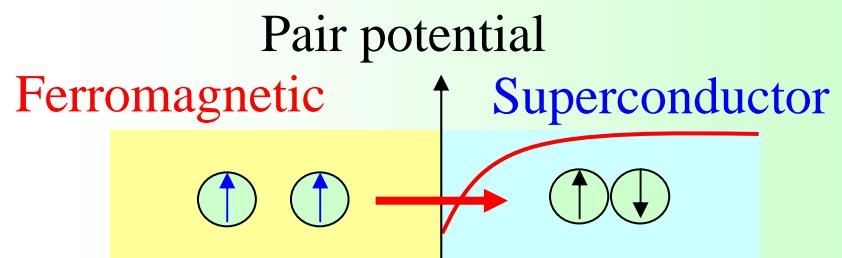
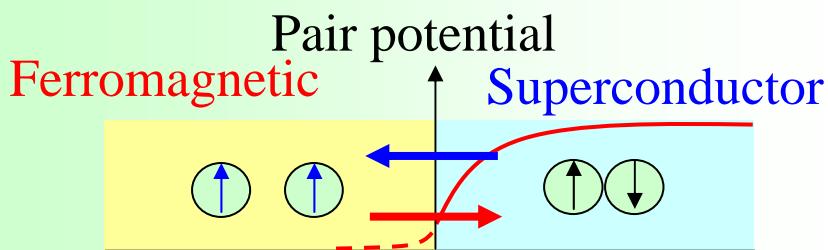


Spin Injection

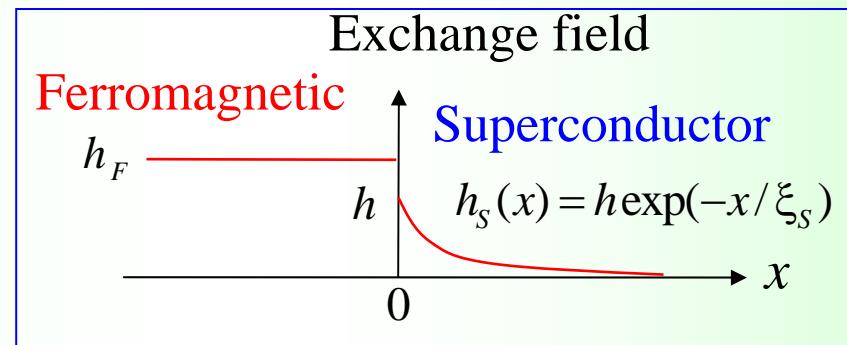


(In,Mn)As / InAs junction

Fe / Si junction



Inverse proximity effect



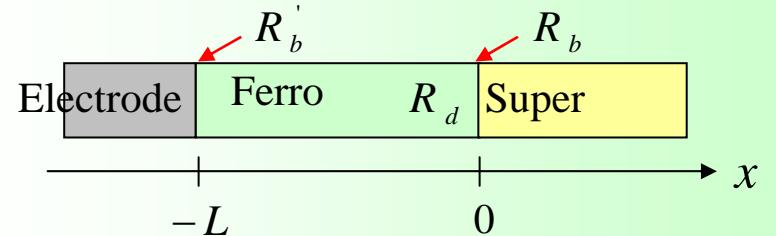
Theoretical Model

Usadel Equation

$$D \frac{\partial}{\partial x} \left(\hat{G} \frac{\partial}{\partial x} \hat{G} \right) + i [\hat{H}, \hat{G}] = 0$$

\hat{G} :Green's function D :Diffusion constant

Mean free path $l \ll L$



Hamiltonian for spins

$$\hat{H} = \begin{cases} (\varepsilon + (-)h_F) \hat{\tau}_3 & -L \leq x \leq 0 \\ [\varepsilon + (-)h_S(x)] \hat{\tau}_3 + i\Delta(x) \hat{\tau}_2 & x > 0 \end{cases}$$

ε : Quasiparticle energy

$\Delta(x)$: pair potential

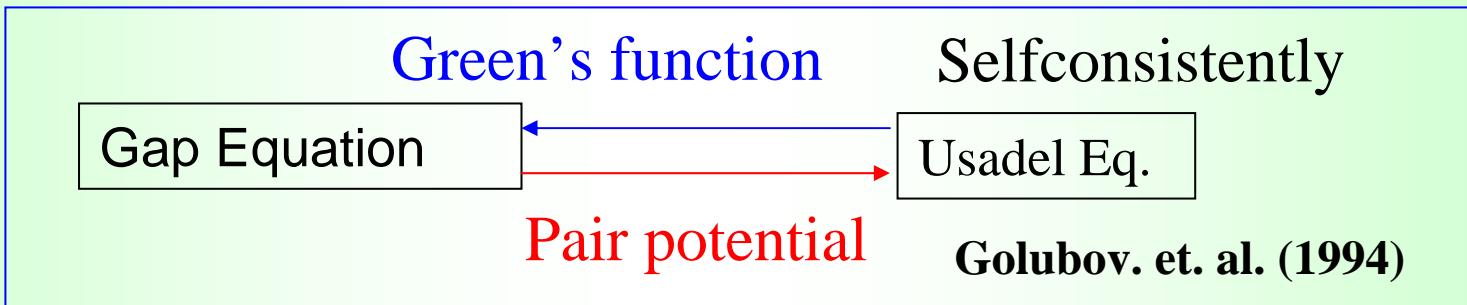
$h_{F(S)}$: Exchange field in F(S)

$\hat{\tau}_i$: Pauri matrix ($i = 1, 2, 3$)

Boundary condition: Conservation at the interface

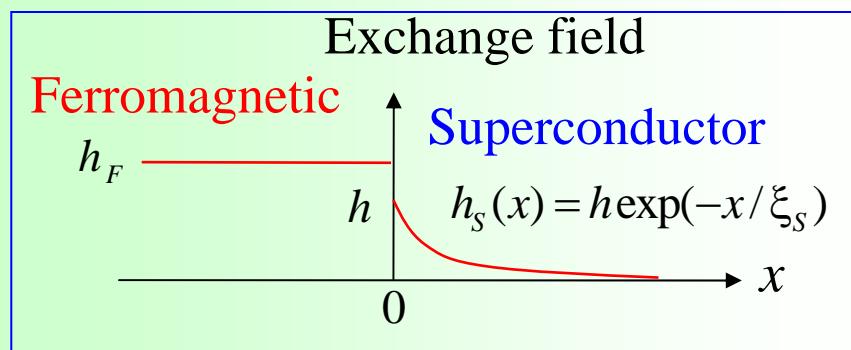
$\Delta(x \rightarrow \infty) = \Delta_0$: Bulk value

Calculation Method



Green's function
↓

DOS, Conductance, etc.



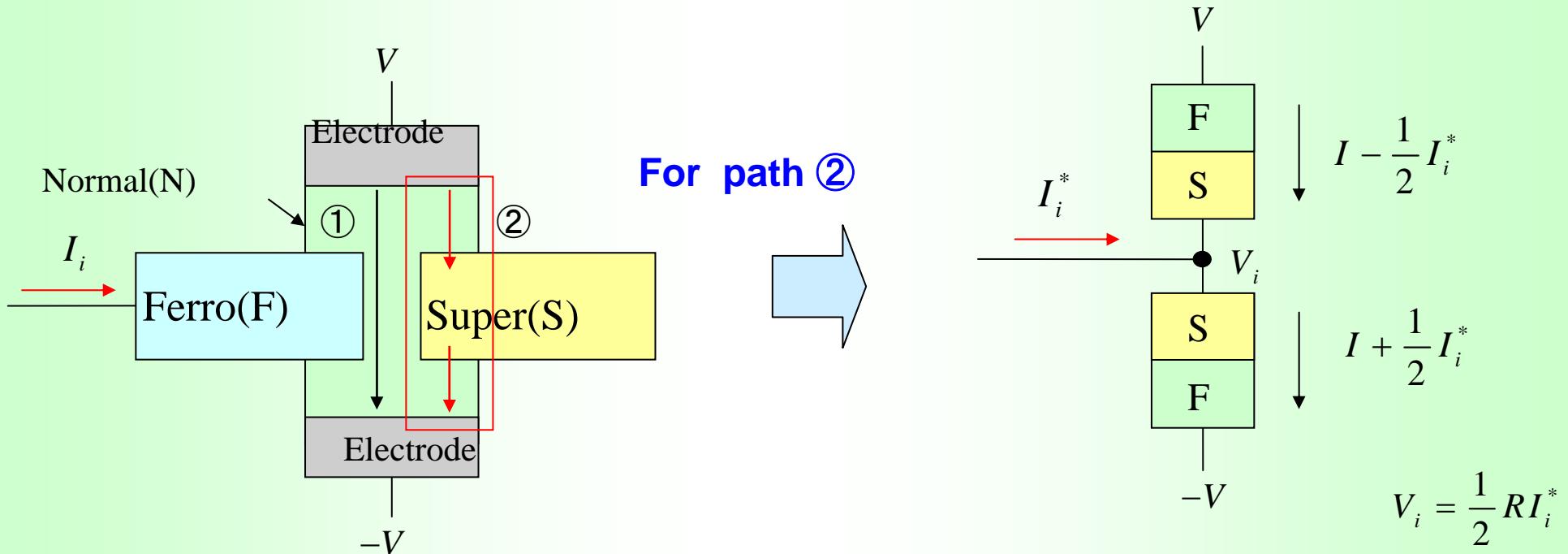
Constant h_F

Exponential decay of h_s

$h_F \geq h$ Due to the barrier at the interface

$\xi_{F(S)}$: F(S) Coherence length

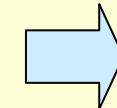
Model for the Experiment



V_i Virtual voltage for the current injection

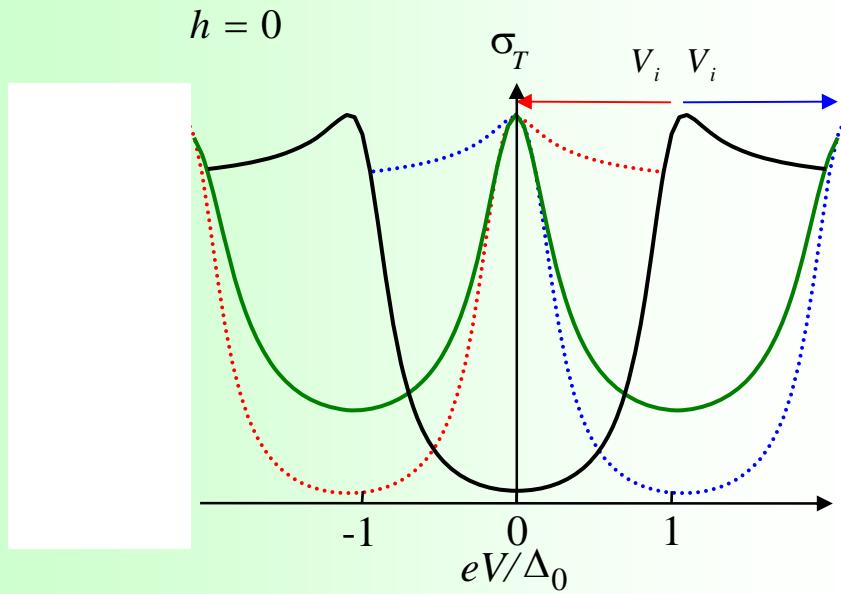
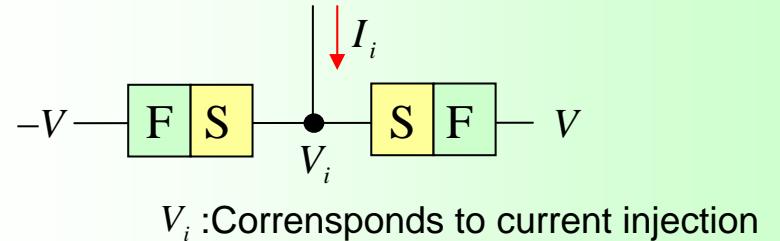
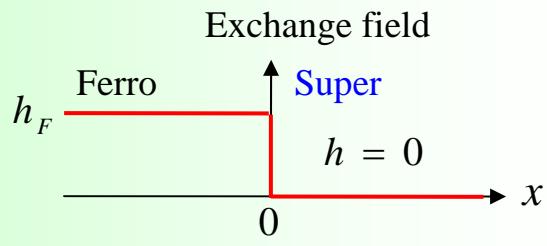
Current injection from Ferromagnetic

Normal

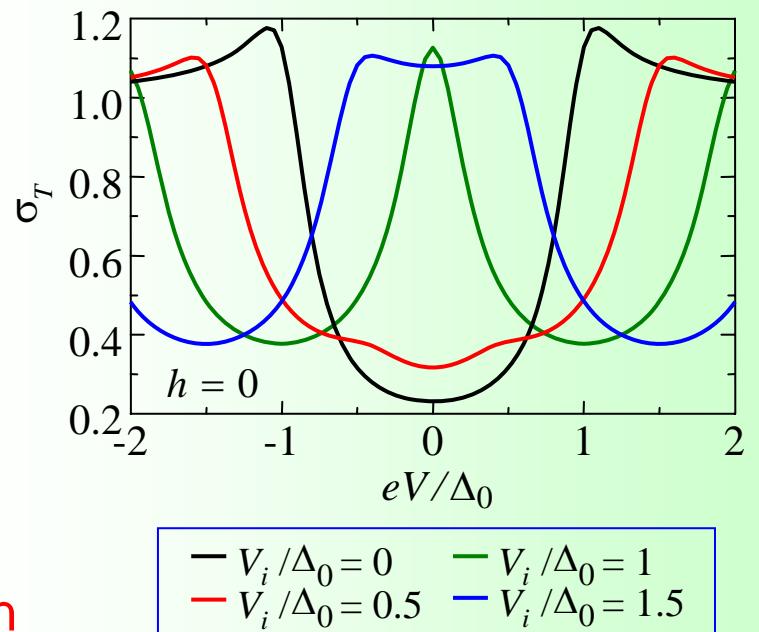


Ferromagnetic

Conductance for a FS-SF junction (1)

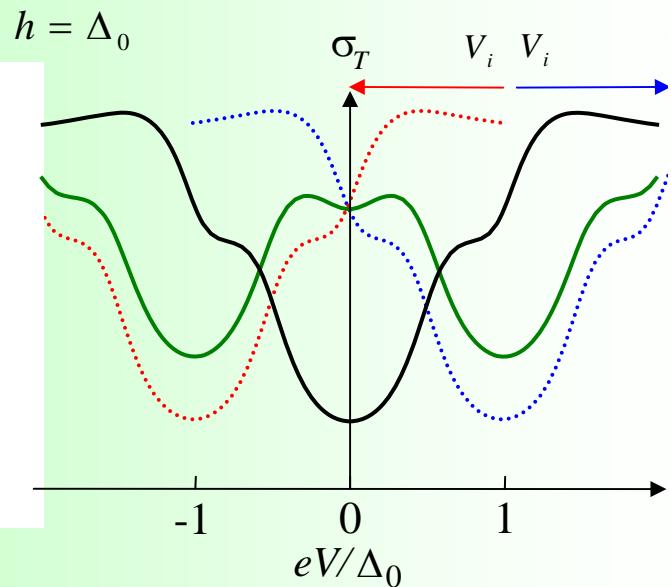
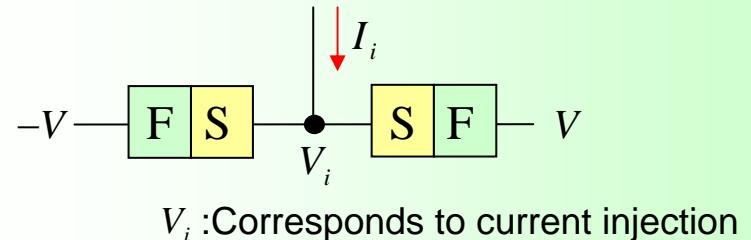
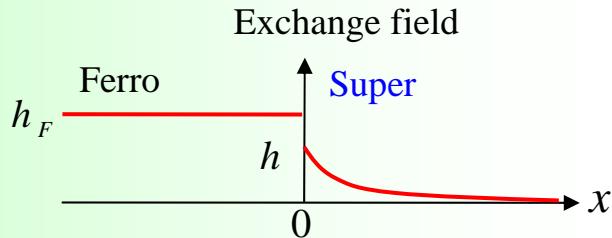


Zero-bias conductance peak by current injection

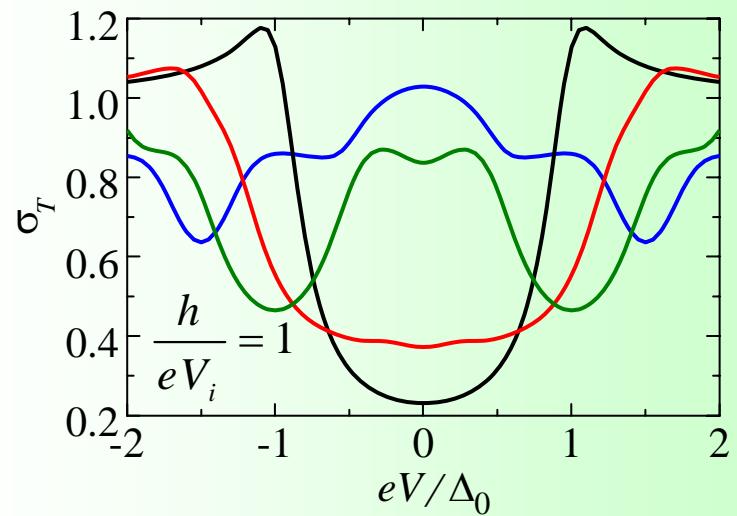


Conductance for a FS-SF junction (2)

Inverse proximity effect

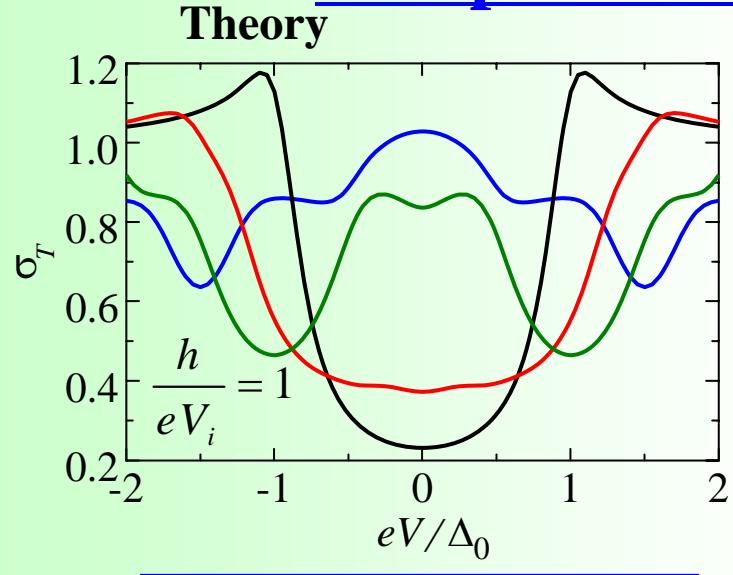


Zero bias peak is suppressed by the synergistic effect of the current injection and exchange field in the S.

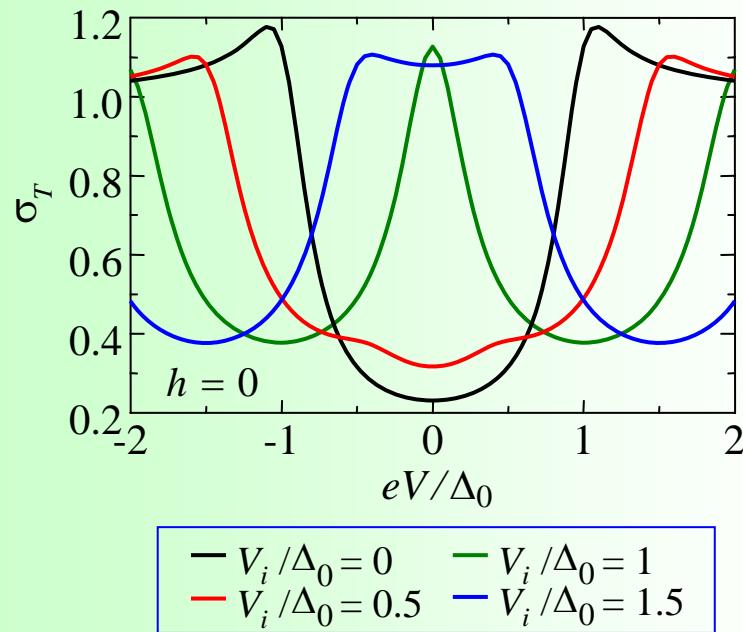


$V_i/\Delta_0 = 0$	$V_i/\Delta_0 = 1$
$V_i/\Delta_0 = 0.5$	$V_i/\Delta_0 = 1.5$

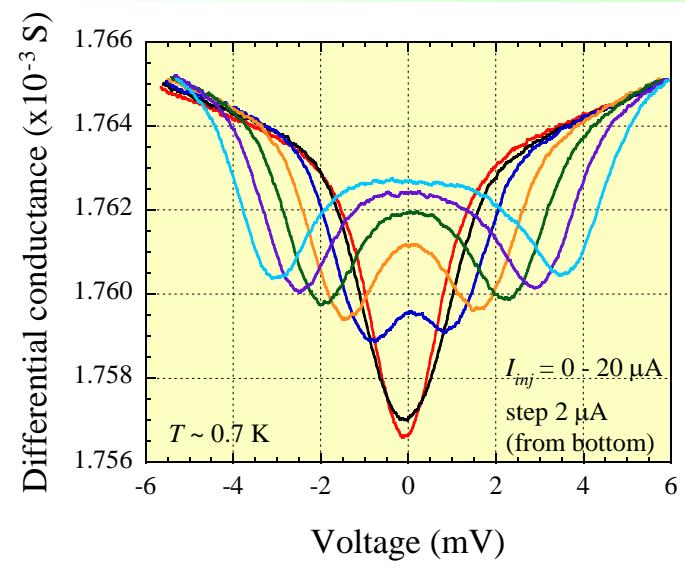
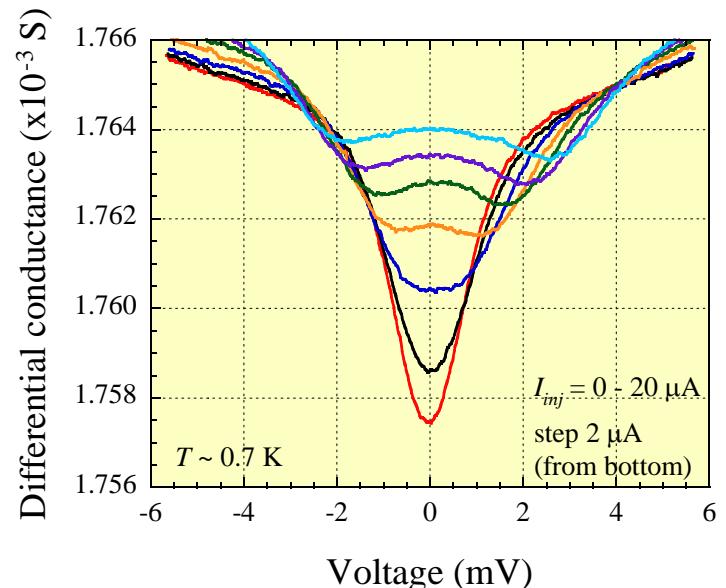
Comparison with experimental results



$$h_S \neq 0$$



$$h_S = 0$$



Summary

1. Nb/**p-InMnAs**/Nb junctions.

- Suppression of Andreev reflection due to spin polarization in p-InMnAs

2.Nb/**n-InAs**/ferromagnetic **p-InMnAs** junction

- We can study the conductance of two types of junctions; one is with the **inverse proximity effect** and the other is without the inverse effect.
- Our theoretical model explains both experimental results.