

# Interactions between Domain Walls and Spin-polarized Currents

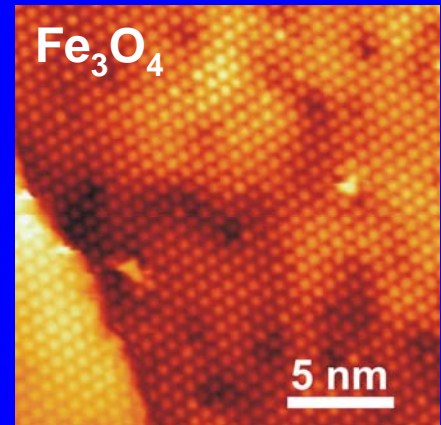
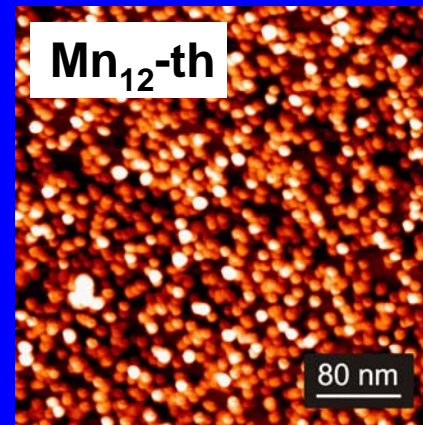
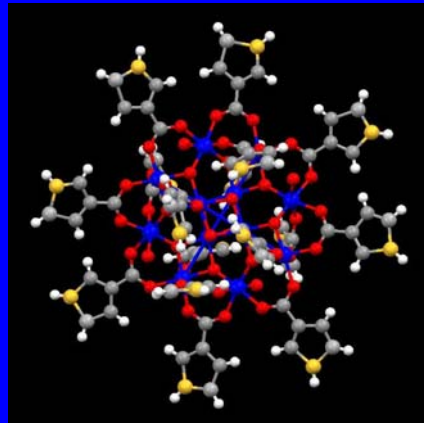
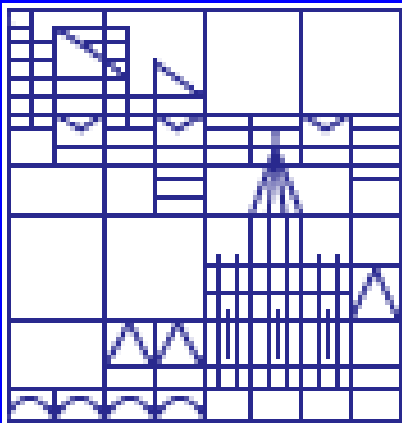
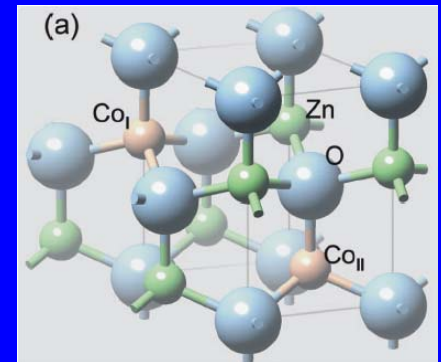
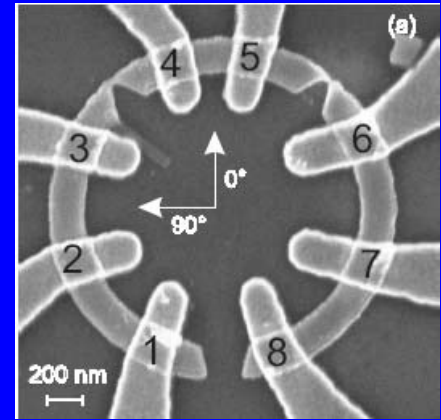
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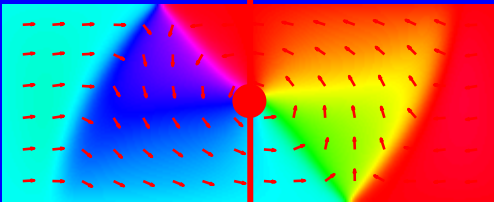
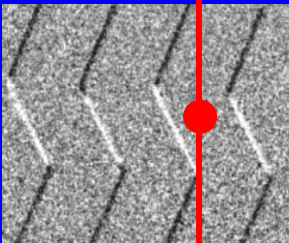
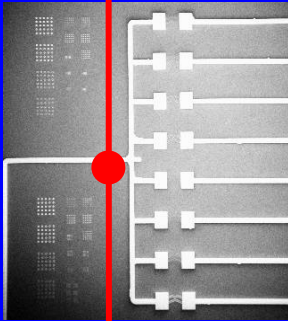


# Present Research Topics

- Spin-dependent transport phenomena
- Interaction of spin-polarized currents with domain walls
- Micromagnetic simulations
- Halfmetallic ferromagnets (HMF)
- Diluted magnetic oxidic semiconductors
- Single molecule magnets (SMM)



# Outline



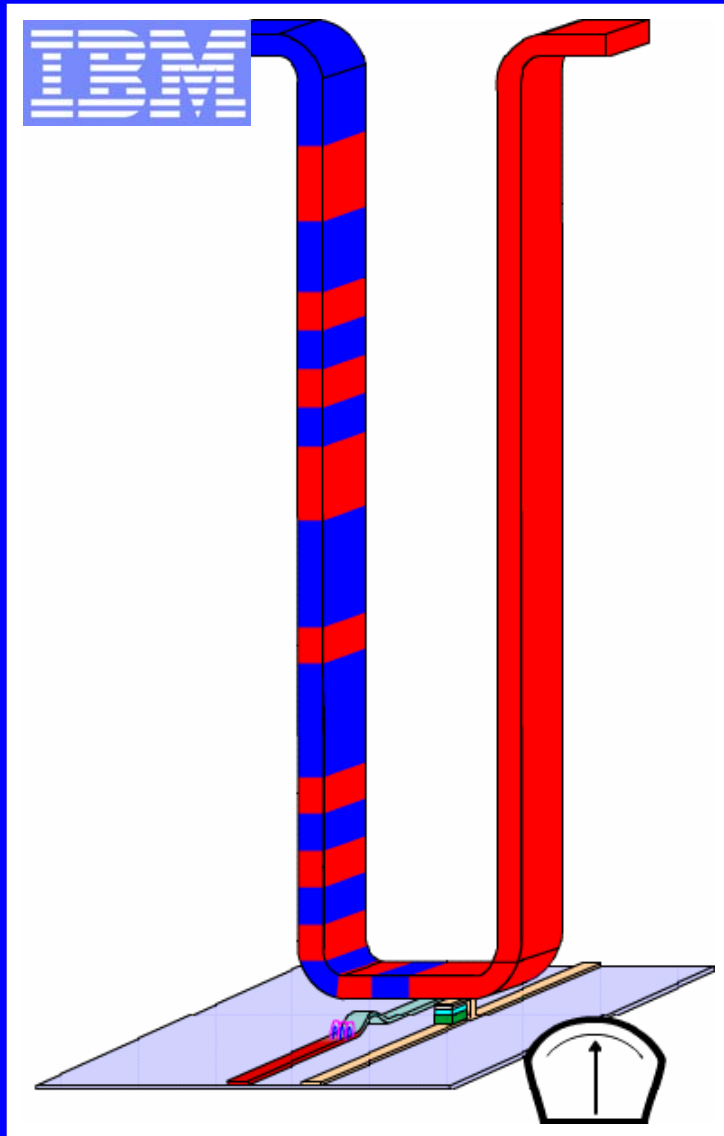
- Motivation (race track memory)
- Resistivity contributions due to domain walls (AMR and DWMR)
- Current-induced domain wall propagation (CIDP): an overview
- Direct observation of CIDP in magnetic zig-zag lines
- Current-induced DW transformations
- The role of temperature

# Research Group/Collaborations

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- *G. Faini (LPN CNRS Marcoussis)*
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- *A.D. Kent (New York University)*

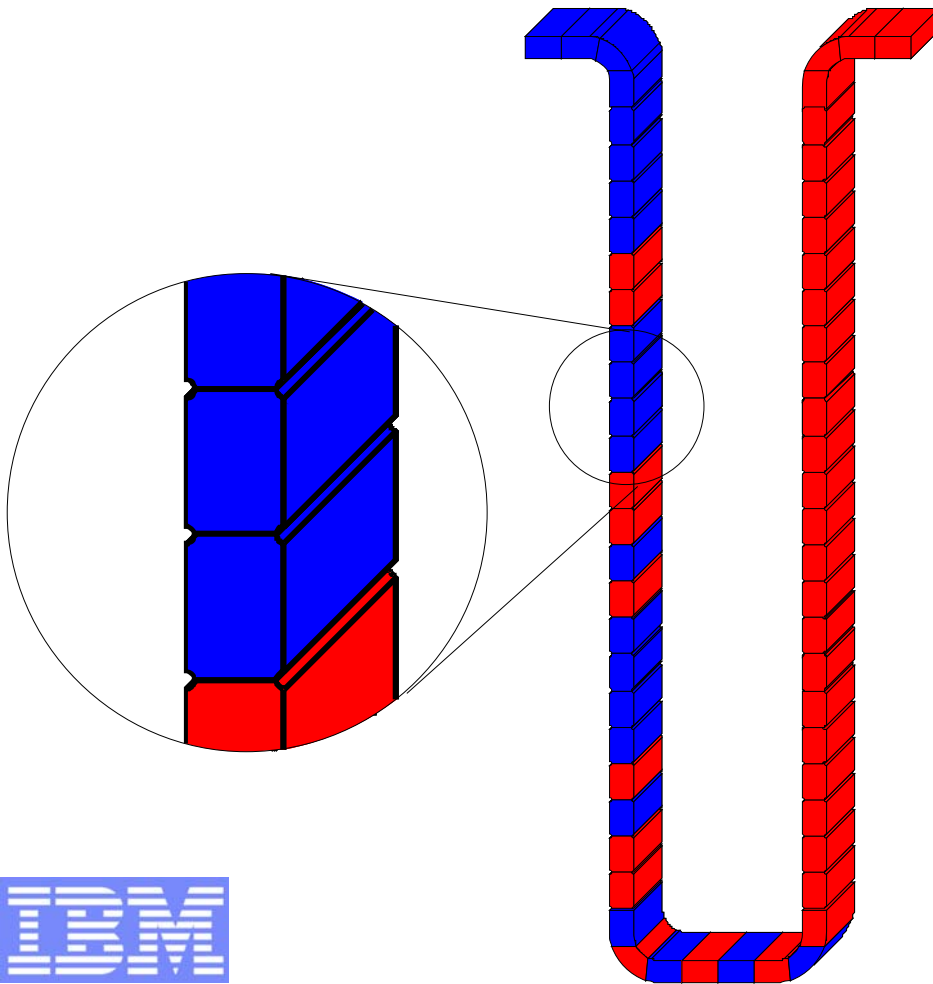
# Race Track Memory: Open up the 3<sup>rd</sup> Dimension!



- Magnetic domains represent the bits
- Approximately 100 Bits per cell (cell width 100 nm)
- Domain wall positioning by current-induced domain wall motion
- Read: magnetoresistive read elements
- Write: local Oersted fields



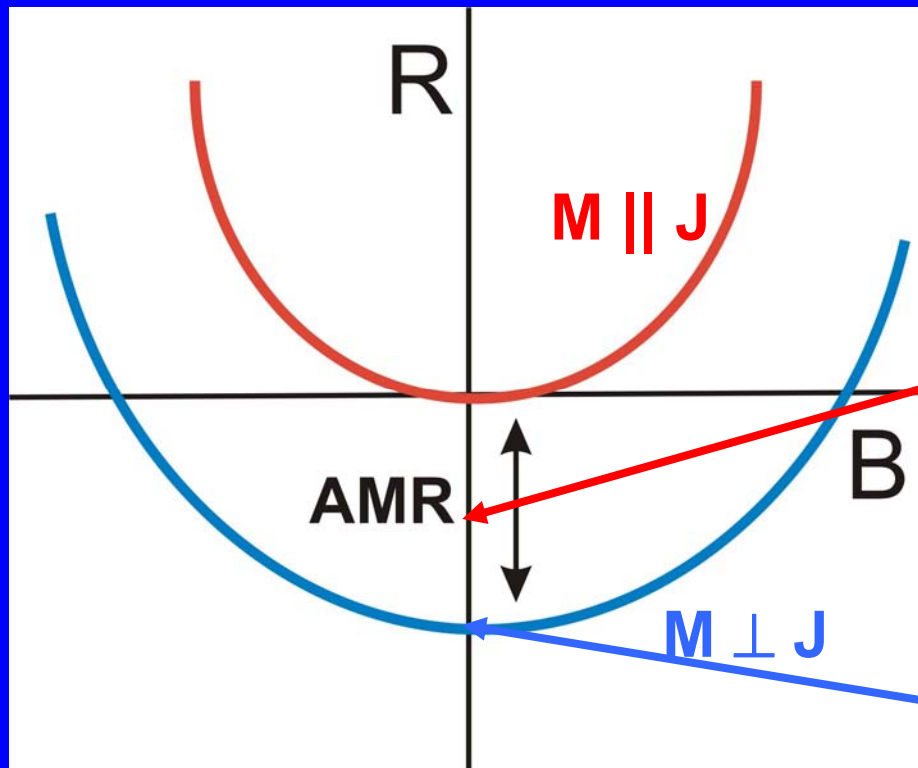
# Race Track Memory: Requirements



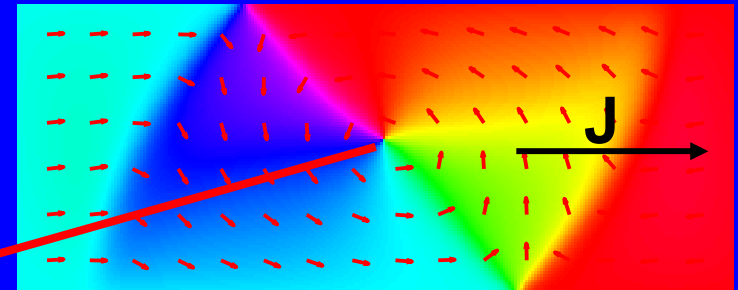
- Domains and domain walls, which can be tailored (spin structure, etc.)
- Well-defined domain wall positions
- We need to select the wall motion direction and move all the domain walls synchronously → ***current-induced motion***
- Reproducibility

# **Magnetoresistive Effects in Presence of Domain Walls**

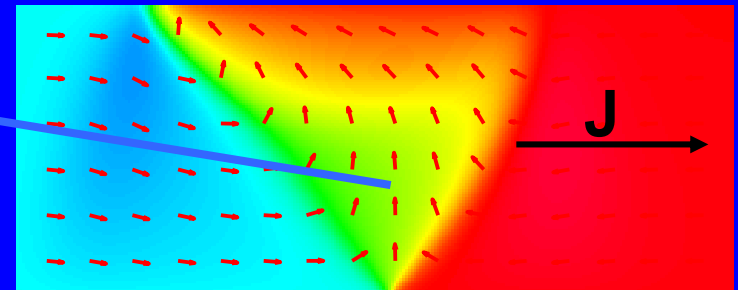
# Anisotropic Magnetoresistance Contribution (AMR)



Vortex DW (thick, wide wires)



Transverse DW (thin, narrow wires)

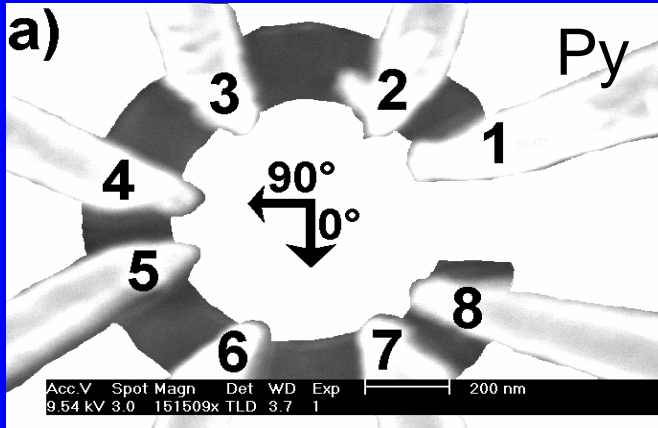


- Origin of AMR: spin-orbit-coupling ( $J$  vs  $M$ )
- MR in the %-range
- Can be used to determine the location of a DW

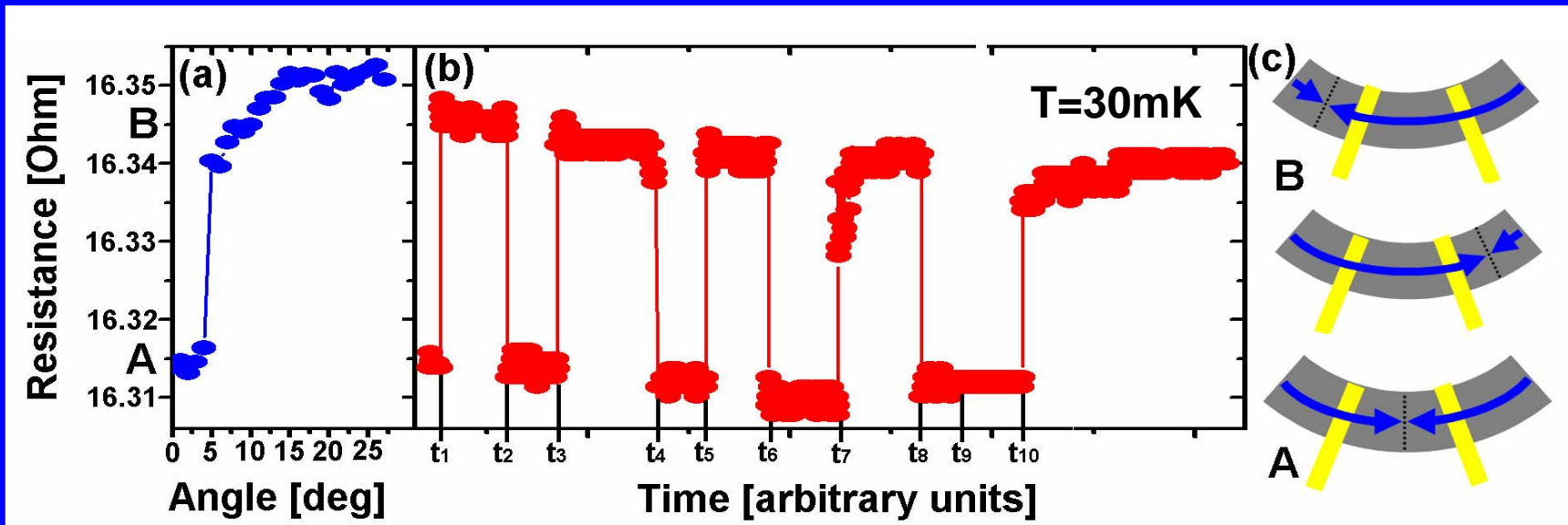
*PRL* **80**, 5639 (1998)  
*APL* **86**, 032504 (2005)



# Magnetoresistive Observation of CIDP in Magnetic Rings

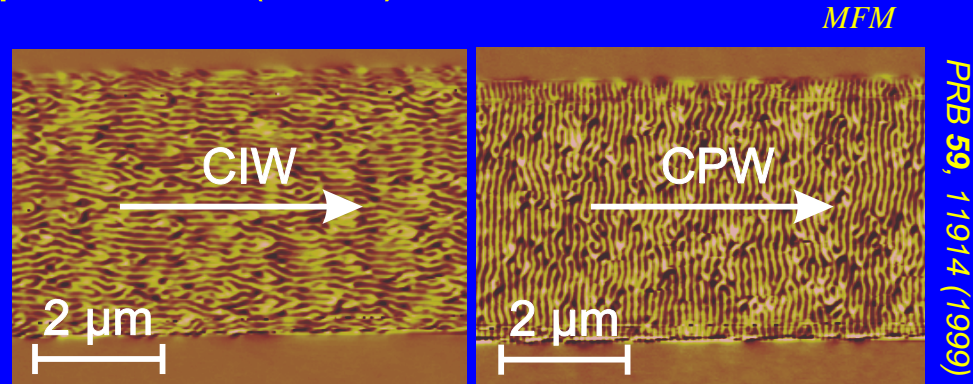


- Voltage measurement between 6 and 7
- Lock-in current at 1
- Current pulses at 2; ground at 8
- Level A → DW between contacts
- Level B → DW outside contacts
- DW can be reversibly moved between positions A and B by current pulses with opposite polarity ( $20 \mu\text{sec}$ ;  $2 \times 10^{12} \text{ A/m}^2$ )



# Domain Wall Magnetoresistance (DWMR)

- **Current-In-Wall (CIW)** and **Current-Perpendicular-To-Wall (CPW)** geometry of epitaxial Co(0001) wires:



- Spin dependent scattering in the presence of **domain walls** leads to an additional resistivity contribution (P.M. Levy et al., PRL **79**, 5110 (1997)):
- Use ferromagnets with a **large uniaxial anisotropy**:

$$a = \pi \sqrt{\frac{A}{K_U}}$$

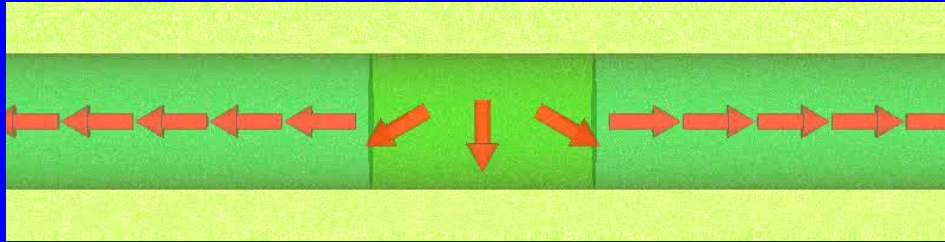
$a_{\text{Fe}} = 40 \text{ nm}$   
 $a_{\text{Co}} = 15 \text{ nm}$   
 $a_{\text{FePt}} < 5 \text{ nm}$

- **Estimation:**  $\text{DWMR}_{\text{CPW}}(\text{Fe, Co, FePt}): < 1\%, \sim 2\%, > 10\%$

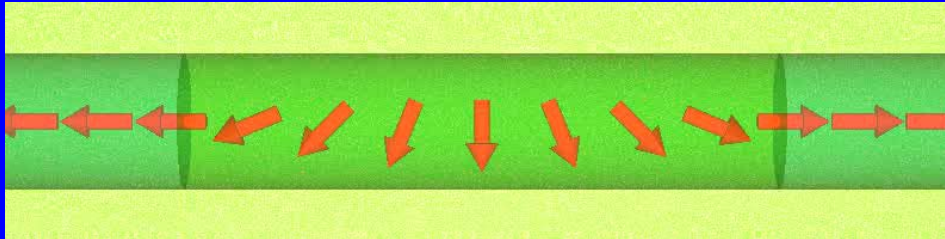
# **Current-induced Domain Wall Propagation in Magnetic Nanostructures**

# Current-induced Domain Wall Propagation (CIDP)

a) Narrow Wall: Momentum transfer to DW

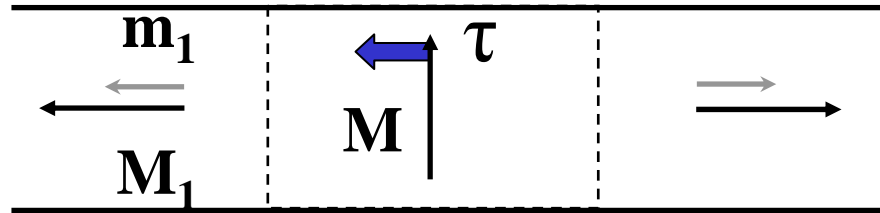


b) Wide wall: Angular momentum transfer



*G. Tatara et al., PRL* **92**, 86601 (2004).

# Current-induced Domain Wall Propagation (CIDP)



- Magnetization dynamics: implicit Landau-Lifshitz-Gilbert equation

$$\dot{\vec{m}} = \gamma_0 \vec{H} \times \vec{m} + \alpha \vec{m} \times \dot{\vec{m}}$$

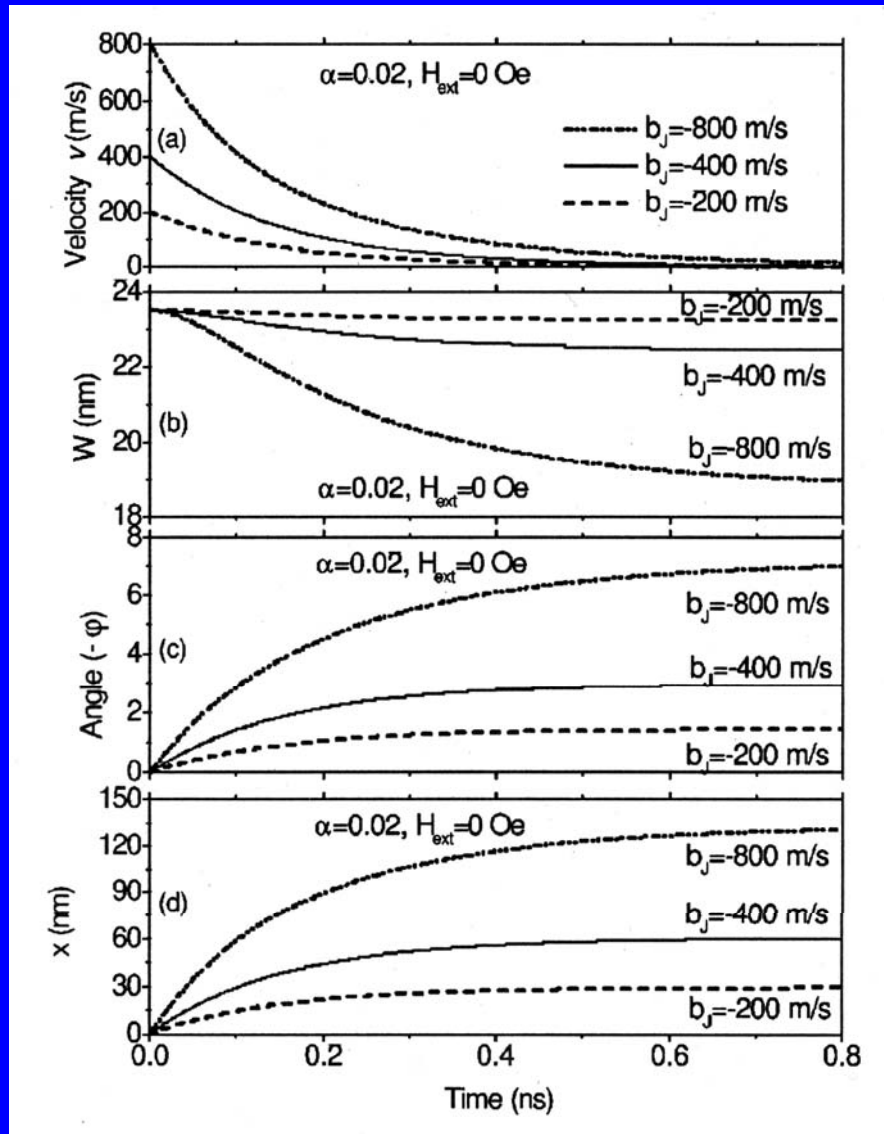
- Spin-transfer model:

$$\dot{\vec{m}} = \gamma_0 \vec{H} \times \vec{m} + \alpha \vec{m} \times \dot{\vec{m}} - (\vec{u} \cdot \vec{\nabla}) \vec{m} - \beta \vec{m} \times [(\vec{u} \cdot \vec{\nabla}) \vec{m}]$$

$$\vec{u} = \vec{j} g P \mu_B / (2e M_s), \quad \beta = (\lambda_J / \lambda_{sf})^2$$

- Angular momentum conservation → „spin transfer“.
- Domain walls move in the direction of the electron flow.
- The effect is proportional to the current density  $j$  and the spin-polarization  $P$  (and inversely to  $M_s$ ).

# Theoretical Models (Pure Adiabatic Processes: $\beta=0$ )



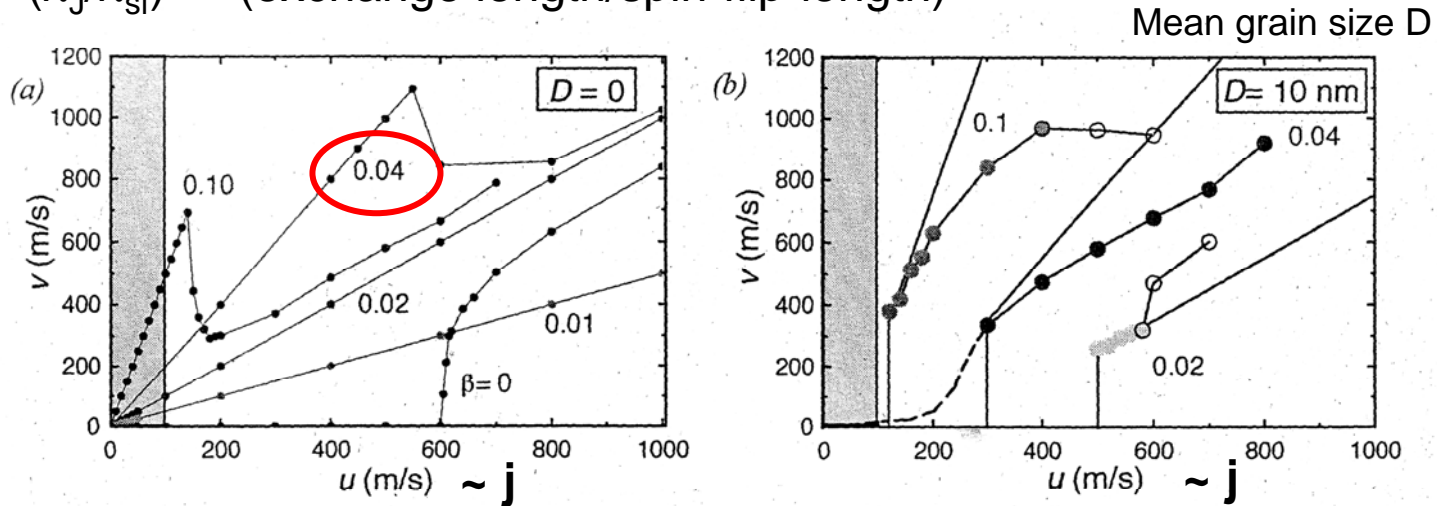
- **Assumption:** adiabatic process, i.e. magnetic moment of conduction electrons is parallel to the local magnetization.
- **Adiabatic** spin-transfer torque on the magnetization.
- DW has **maximum velocity** at the initial application of the current.
- DW **velocity decreases to zero** as the DW begins to deform during motion ( $W$ : DW width).
- DW is **unable to maintain** the wall movement.





# Theoretical Models (Non-adiabatic Processes: $\beta \neq 0$ )

$$\beta = (\lambda_J / \lambda_{sf})^2 = (\text{exchange length} / \text{spin-flip-length})^2$$

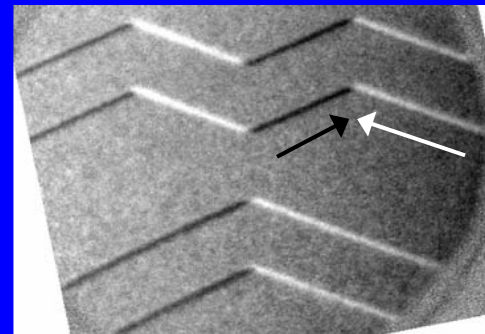
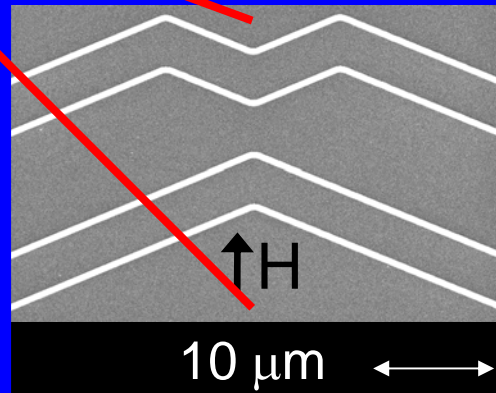
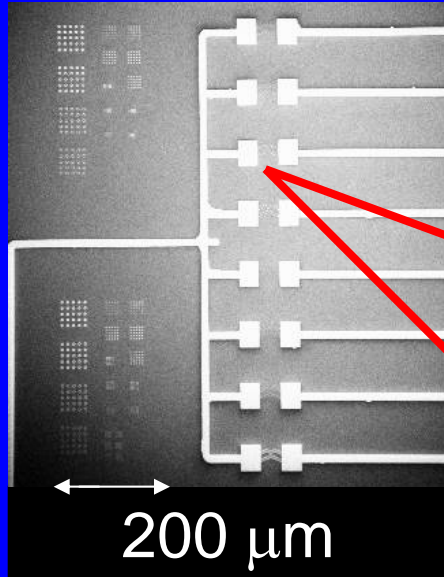


A. Thiaville et al., EPL **69**, 990 (2005).

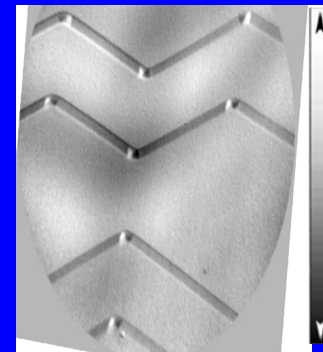
- **Corrections** to perfect adiabaticity and pure local spin transfer, meaning a modification of the initial spin transfer torque by a second order quantity.
- **For  $\beta=0$ :** absence of DW motion for  $u < u_c$ .
- **For  $\beta \neq 0$ :** DW motion at any finite  $u$ ; DW velocity  $v$  increases with increasing  $\beta$ .
- Valid for transverse and vortex walls.
- Exp. observed threshold currents are much smaller.
- Up to now: **neglect of thermal fluctuations**.

# **Direct Observation of CIDP in Magnetic Zig-zag Lines by XMCD-PEEM Imaging**

# Direct Observation of CDP in Magnetic Zig-zag Lines by XMCD-PEEM Imaging



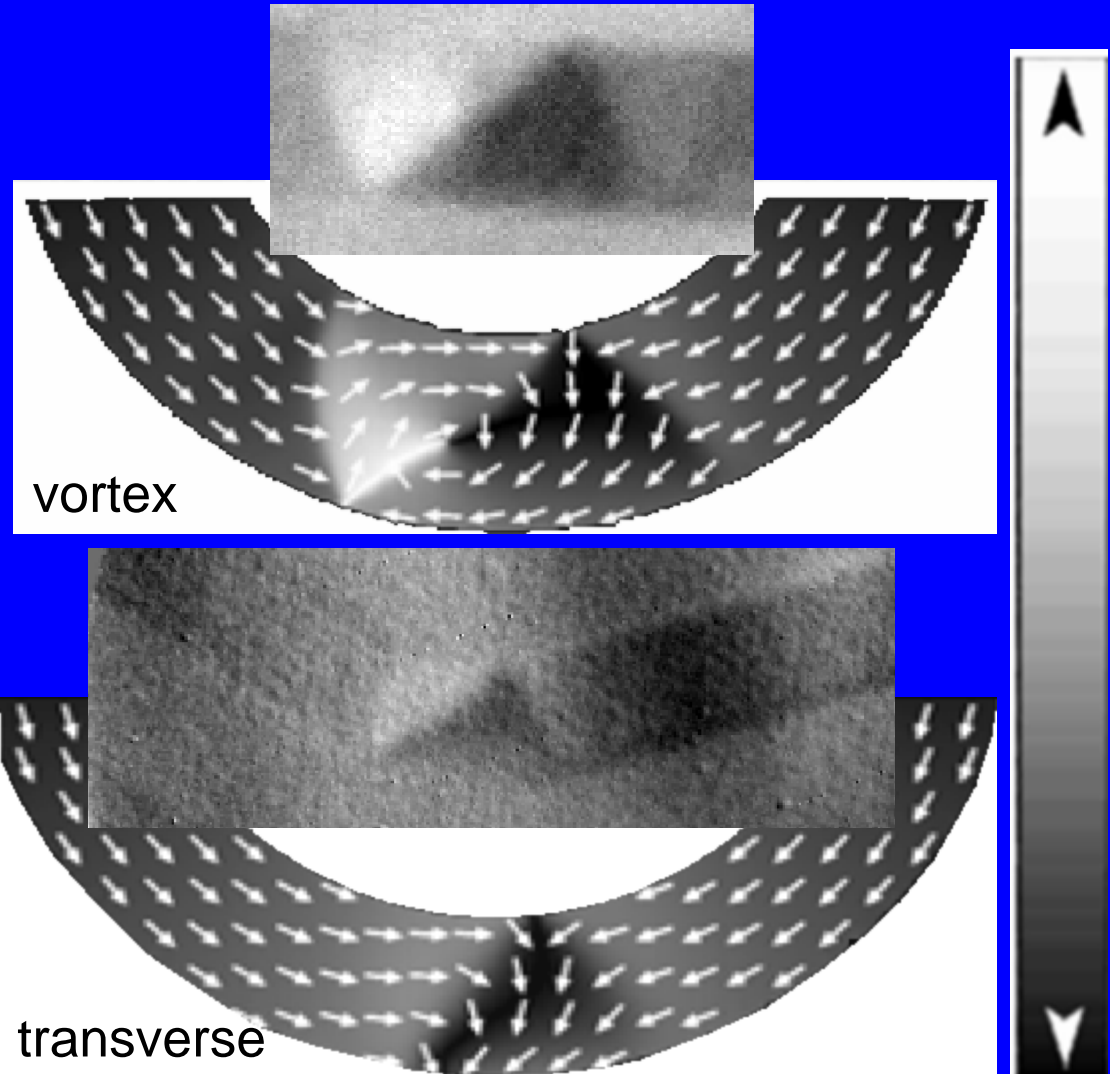
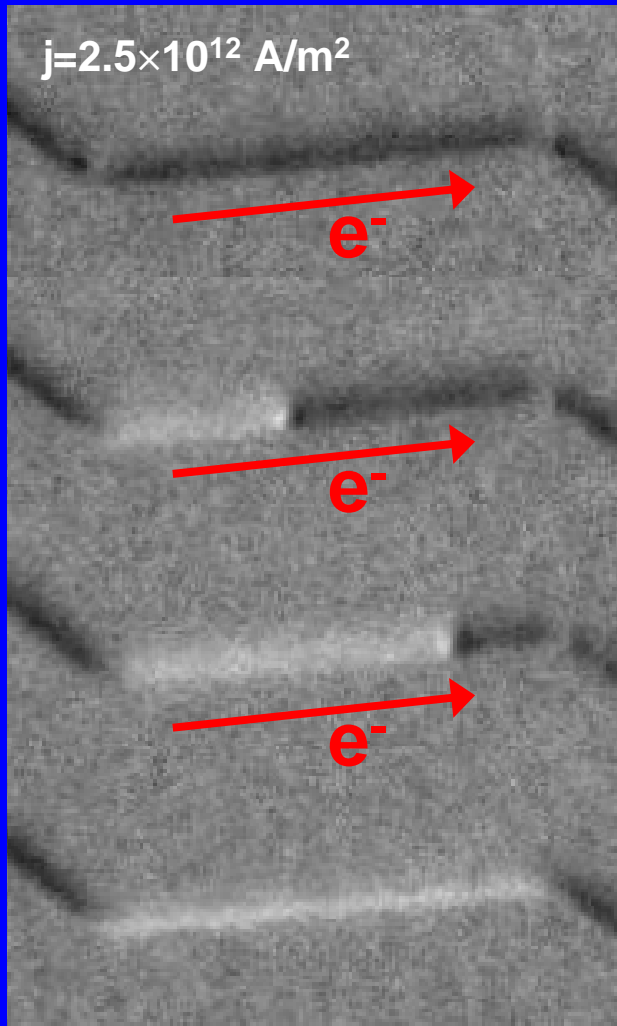
**XMCD-PEEM  
Imaging**



**Vortex Wall**

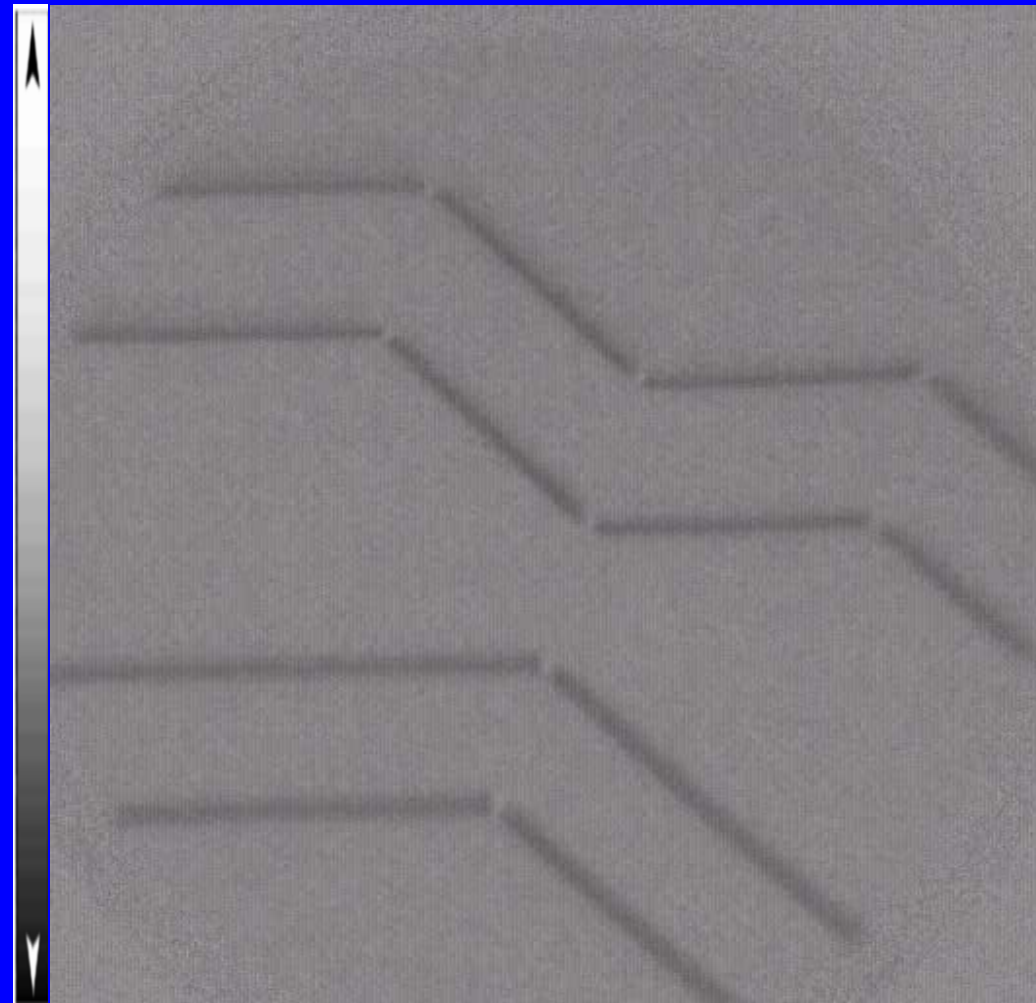
(Zig-zag lines with Au contact pads:  $W=500\text{nm}$ ;  $L=10\mu\text{m}$ ;  $t=10\text{nm}$  Py)

# Direct CIDP Observations with XMCD-PEEM

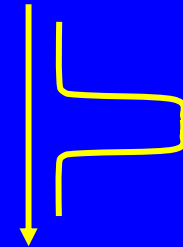
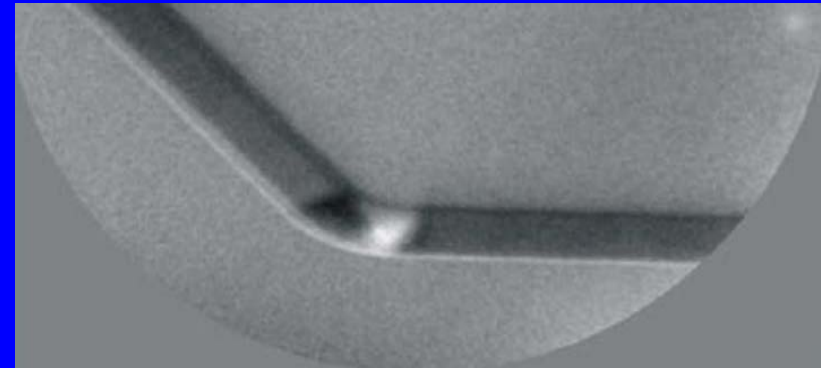


- Domain walls move in the direction of the electron flow.
- High-resolution imaging reveals the domain wall spin structures (vortex, transverse).

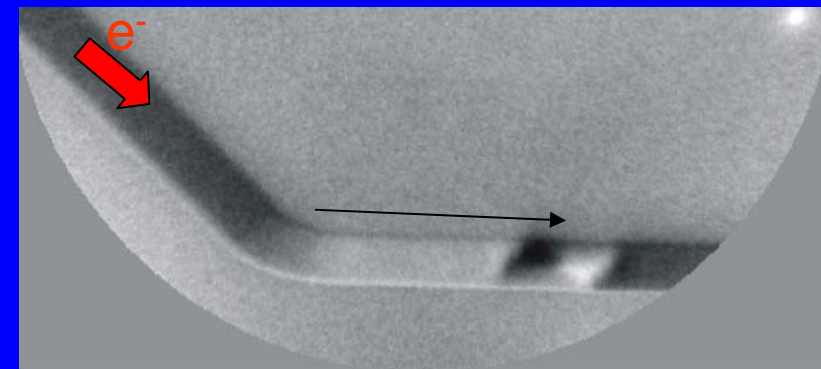
# The Stochastic Nature of CIDP (XMCD-PEEM)



Pulses with 146V-156V, 150 $\mu$ s



Current pulse  
25 $\mu$ s,  $10^{12}$ A/m<sup>2</sup>

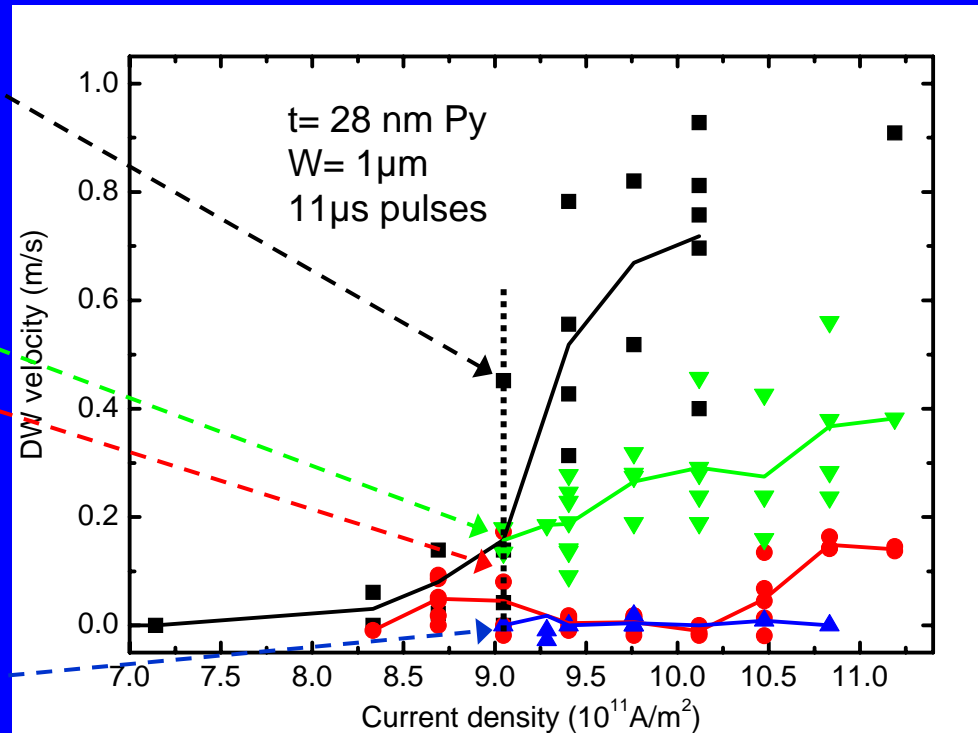
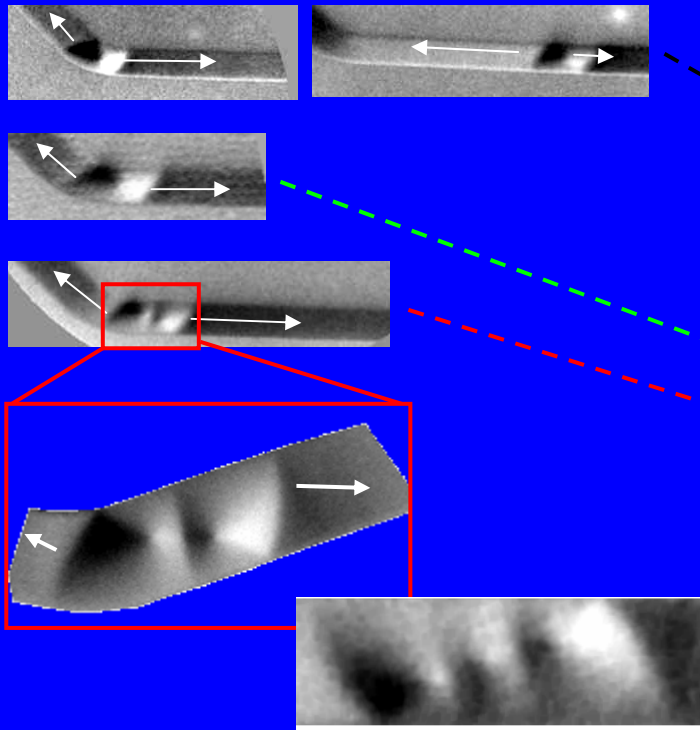


Py; 1 $\mu$ m wide, 28nm thick

# **Current-induced DW Transformations**



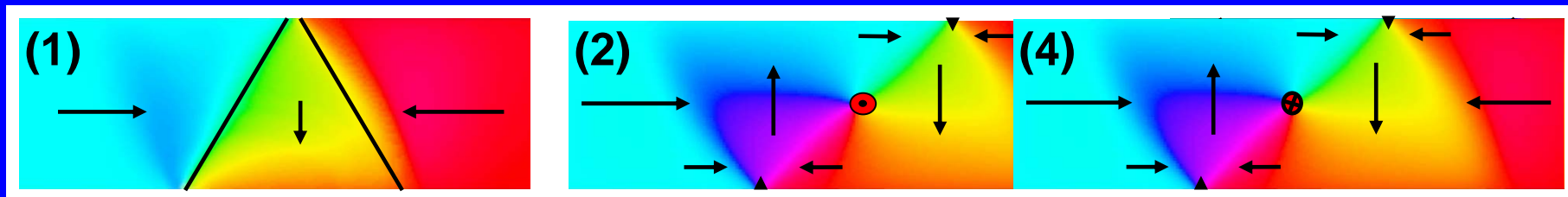
# The Stochastic Nature of CIPD: Vortex Core Nucleation and Annihilation



APL 88, 232507 (2006).

- Velocity of single vortex walls with no transformations **increases with increasing current density** (black squares and black line).
- Velocity depends on the **number of vortices**.
- **Extended vortices** move more slowly (green down triangles).
- **Multi-vortices** (double vortex: red; triple vortex: blue) hardly move.

# Direct Observation of CIDP in „Zig-Zag“ Lines

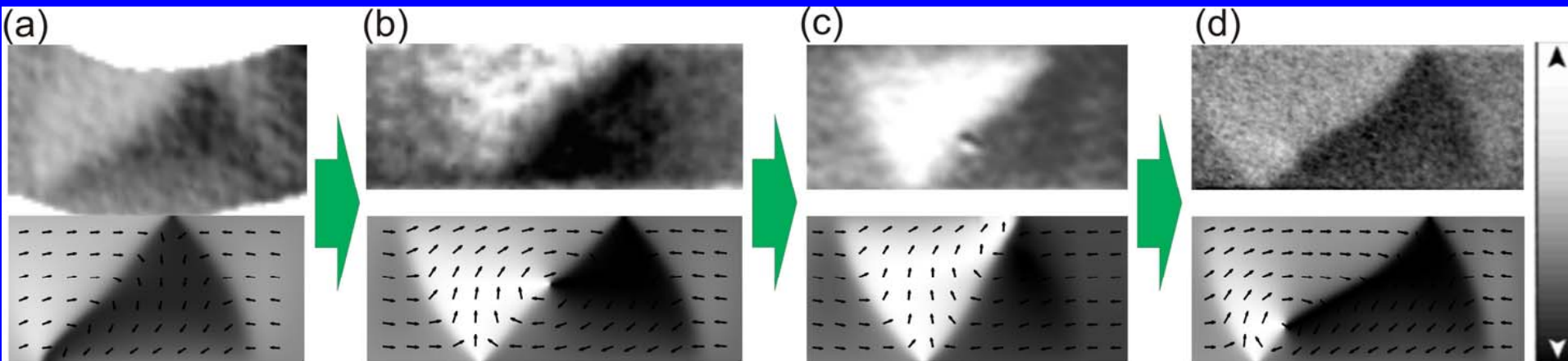


*A. Thiaville et al., EPL 69, 990 (2005).*

- Prediction of periodic transformation of DW type by the nucleation and annihilation of a vortex core: **TW (down) → VW → TW (up) → VW**.
- **TW** is alternating (up/down) but **VW** with the same circulation direction but opposite polarity.

$W = 1.5 \mu\text{m}$ ,  $t = 7 \text{ nm}$ , close to TW-VW phase boundary

*Appears in PRL 2008*



Transverse wall  
(**down**)

Vortex wall (**clockwise**)  
after pulse injection  
( $10^{12} \text{ A/m}^2$ ,  $25 \mu\text{s}$ )

Transverse wall (**up**)  
after pulse injection

Displaced vortex core  
gives direct evidence of  
transformation  
mechanism!

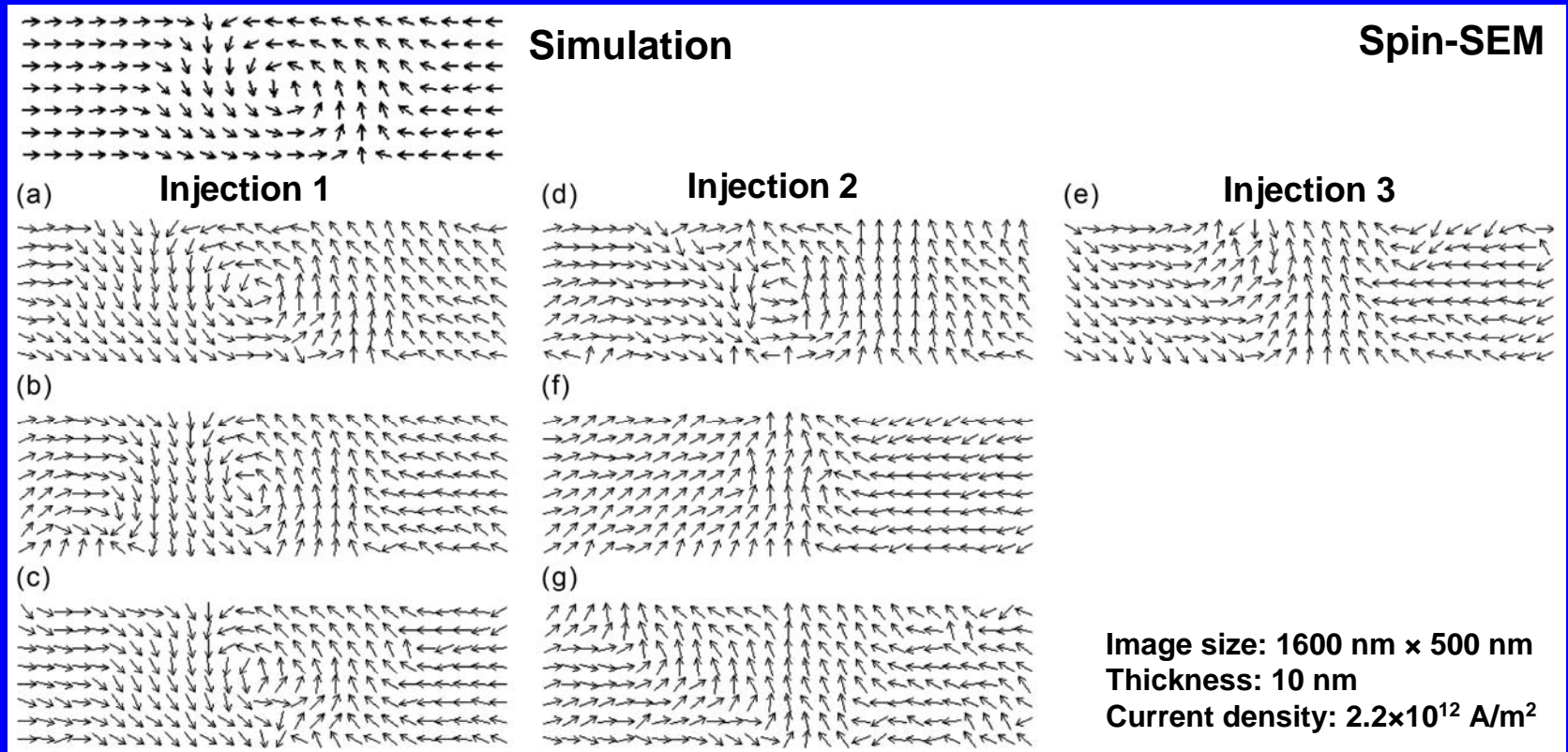
# Direct Observation of CDDP in „Zig-Zag“ Lines

- Vortex core feels a force **perpendicular** to the current; it moves not only in direction of the electron flow (also towards the edges).
- The y-direction movement depends on the polarity of the vortex core; y-velocity is proportional to  $(\alpha - \beta)$  (see: *He et al., PRB 73, 184408 (2006)*).

$$\underline{v} = -\frac{u_x}{1 + \alpha^2} \left( (1 + \beta\alpha) \frac{\partial \underline{S}}{\partial x} + (\alpha - \beta) \underline{S} \times \frac{\partial \underline{S}}{\partial x} \right)$$

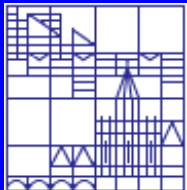
- For large enough currents **the vortex core is expelled** and a TW is formed; then a new VW with **opposite polarity** is nucleated and starts to move to the **opposite edge** of the wire.
- Observation **excludes** a former claim that  $\alpha = \beta$  (PRB 74, 144405 (2006)).
- **Excludes** thermal-activated or defect-induced transformations as these would result in random rotation senses of the VW magnetization.
- **Explains** why in earlier experiments TW stopped for a given current density but vortex walls move (pinning at edge irregularities stronger).

# DW Spin Structure vs Number of Pulse Injections



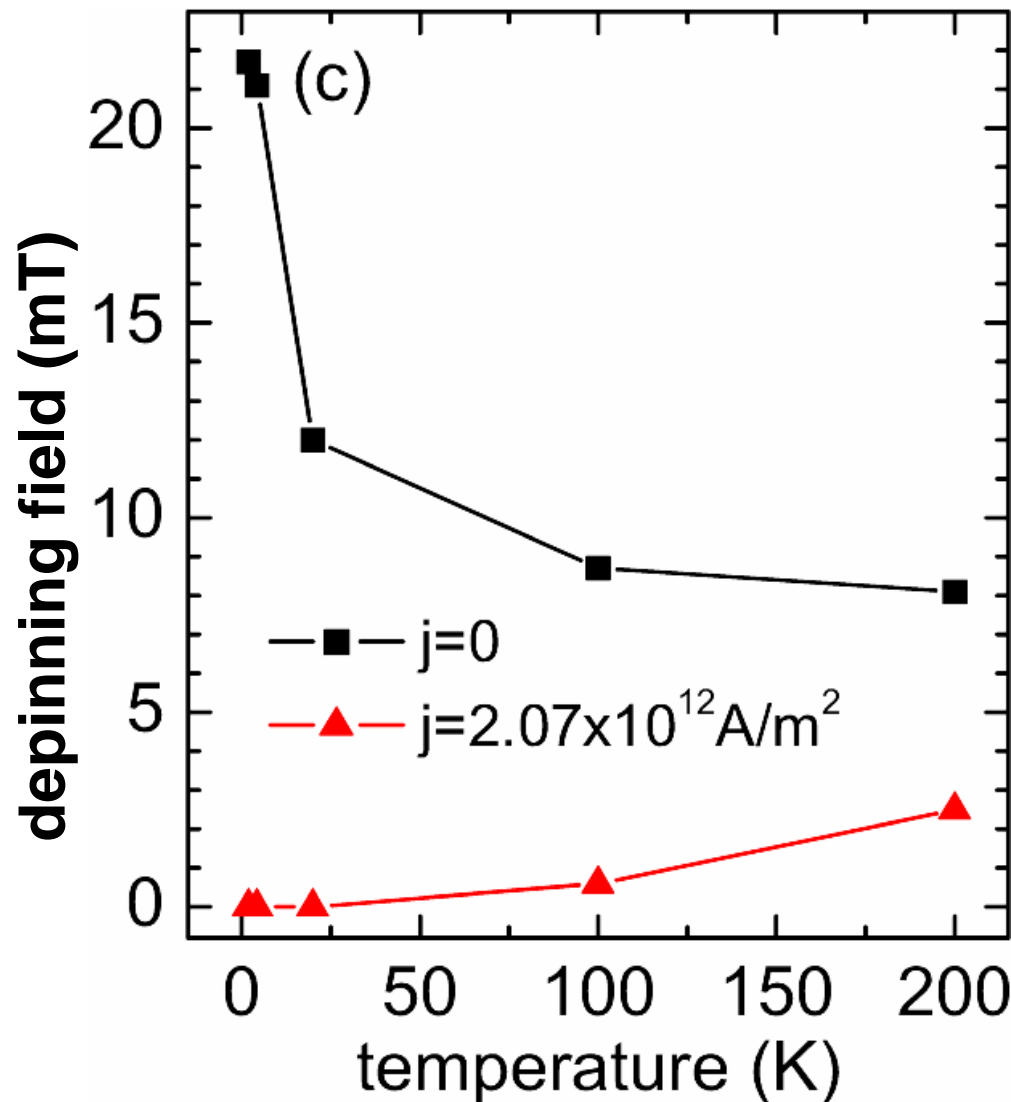
PRL 95, 026601 (2005).

- After the first current injection all walls are of vortex-type.
- After a few injections all three walls have stopped moving and undergone a drastic transformation to a distorted transverse wall.



# **Role of Temperature**

# Role of Temperature

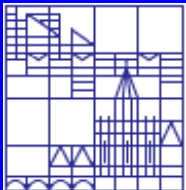


- The magnetic field needed to depin a domain wall **decreases** with **increasing** current density.
- At 0 current, the depinning fields **decrease** with **increasing** temperature, at higher currents the opposite occurs.
- Spin torque effect is **more efficient** at low temperatures!
- Possible explanation of discrepancies between 300K observations and 0K calculations: asymmetric **generation of spin waves**.

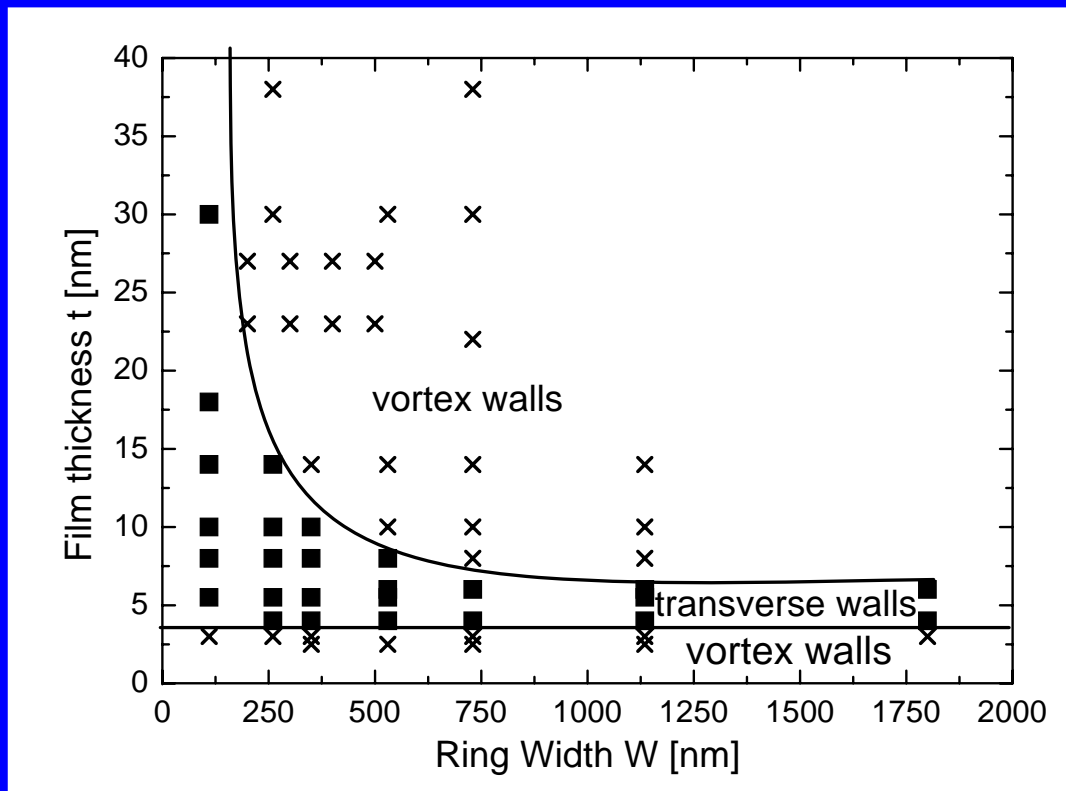


# Summary

- Clear observation of current-induced domain wall propagation (CIDP).
- Modification of the domain wall spin structure by spin-polarized currents (stochastic process).
- VW-TW transformations by current pulses; good agreement with micromagnetic simulations.
- Critical current density for CIDP increases with increasing temperature (spin wave generation?).



# Domain Wall Phase Diagram for Permalloy Rings



*APL* **88**, 52507 (2006).